University of London

## EXAMINATION FOR INTERNAL STUDENTS

## For The Following Qualifications:-

M.Sc. $\quad P G D i p$

Frontiers in Mathematical Modelling and its Applications

COURSE CODE : MATHGM05

DATE
: 15-MAY-06

TIME
$: 10.00$

TIME ALLOWED : 2 Hours

All questions may be attempted but only marks obtained on the best four solutions will count.
The use of an electronic calculator is not permitted in this examination.

1. (a) You are given a sample of a linear viscoelastic material and an oscillatory rheometer. Describe briefly how you would measure the linear rheology of the material.
(b) Given the definition

$$
G^{*}(\omega)=i \omega \int_{0}^{\infty} G(s) \exp [-i \omega s] \mathrm{d} s
$$

of the complex shear modulus, give expressions for the storage modulus $G^{\prime}$ and loss modulus $G^{\prime \prime}$ of the material.
(c) Prove for the experiment you described in (a) that the shear stress is

$$
\sigma=G^{\prime} \gamma(t)+\frac{G^{\prime \prime}}{\omega} \dot{\gamma}(t)
$$

where $\gamma$ is the shear displacement and $\dot{\gamma}$ the shear rate.
(d) Calculate the storage and loss moduli for a material with a single exponential relaxation time, $\tau$. For what value of $\omega$ are the two moduli equal?
2. (a) A vortex of strength $\kappa$ is present at the point $z_{0}$ in a region $R_{z}$ of the complex $z$-plane so that the complex potential $w(z)$ is analytic everywhere in $R_{z}$ except at the point $z_{0}$, with

$$
w(z) \rightarrow-(\mathrm{i} \kappa / 2 \pi) \log \left(z-z_{0}\right), \quad \text { as } \quad z \rightarrow z_{0} .
$$

The associated complex velocity is

$$
u-\mathrm{i} v=\frac{\mathrm{d} w}{\mathrm{~d} z}
$$

The region $R_{z}$ is mapped into the region $R_{\zeta}$ of the complex $\zeta$-plane by the conformal mapping, analytic within $R_{z}$, given by

$$
\zeta=f(z), \quad z=f^{-1}(\zeta)=F(\zeta)
$$

(i) Show that the vortex at $z_{0}$ maps to a vortex at $\zeta_{0}=f\left(z_{0}\right)$ and find its strength.
(ii) Find an expression for the fluid velocity at a point $\zeta \neq \zeta_{0}$ in $R_{\zeta}$ in terms of $w(z)$.
(iii) Find an expression for the velocity in $R_{\zeta}$ of the vortex at $\zeta_{0}$ in terms of $w(z)$, given that there is no self-induced motion of a point vortex.
(b) Suppose that there are point vortices of strength $\kappa_{1}, \kappa_{2}, \ldots, \kappa_{n}$ with instantaneous positions $\zeta_{1}, \zeta_{2}, \ldots, \zeta_{n}$ moving in infinite space and that the vorticity is zero elsewhere. Then the quantity

$$
H\left(\zeta_{1}, \zeta_{2}, \ldots, \zeta_{n}\right)=-(1 / 4 \pi) \sum_{i=1}^{n} \sum_{j=1,(j \neq i)}^{n} \kappa_{i} \kappa_{j} \log \left|\zeta_{j}-\zeta_{i}\right|
$$

is a constant of the motion.
By considering vortices of strength $\pm \kappa_{0}$ at $\zeta=\zeta_{0}, \bar{\zeta}_{0}$, where - denotes complex conjugate, find $H$ for this vortex pair and hence deduce th , ath of a single vortex in the half-plane $\operatorname{Im} \zeta>0$.
(c) The quantity $H$ remains a conserved quantity under conformal mappings but transforms through

$$
H_{z}\left(z_{1}, z_{2}, \ldots, z_{n}\right)=H_{\zeta}\left(\zeta_{1}, \zeta_{2}, \ldots, \zeta_{n}\right)-(1 / 4 \pi) \sum_{i=1}^{n} \kappa_{i}^{2} \log \left|f^{\prime}\left(z_{i}\right)\right|
$$

for $\zeta=f(z)$.
Use this result with that of (b), and the transformation $\zeta=z^{2}$ to obtain the path of a single vortex in the quarter-plane $\operatorname{Re} z>0, \operatorname{Im} z>0$ and hence sketch it.
3. An equation for the delayed oscillator model of the ENSO cycle is

$$
\begin{equation*}
T_{t}=-\alpha_{1} T-\left(\alpha_{2} / \lambda\right) \tanh [\lambda(T-r T(t-\Delta)] \tag{1}
\end{equation*}
$$

where $\alpha_{1}, \alpha_{2}, \lambda, r$ and $\Delta$ are positive constants, with $r<1$ and $\alpha_{2}>\alpha_{1}$, and $T(t)$ represents the sea surface temperature anomaly in an equatorial region of a bounded ocean.

Describe briefly the physical process that is responsible for the delay time $\Delta$.
Show that when $\alpha_{2}-\alpha_{1}>\alpha_{2} r$ there are at least three different equilibrium values for $T$.

With suitable scaling, an approximate non-dimensional form for (1) is

$$
\begin{equation*}
T_{t}=T-T^{3}-\gamma T(t-\delta) \tag{2}
\end{equation*}
$$

where $\gamma=r \alpha_{2} /\left(\alpha_{2}-\alpha_{1}\right)$ and $\delta=\left(\alpha_{2}-\alpha_{1}\right) \Delta$.
Consider the regime with $1 / 2<\gamma<1$, in which an equilibrium value of (2) is $T_{0}=(1-\gamma)^{1 / 2}$.
(a) If $T=T_{0}+T^{\prime}$, where $T^{\prime}$ is a small perturbation, derive the linearised equation

$$
\begin{equation*}
T_{t}^{\prime}=(3 \gamma-2) T^{\prime}-\gamma T^{\prime}(t-\delta) \tag{3}
\end{equation*}
$$

(b) Show that neutral solutions of (3) oscillate with frequency $\sigma_{I}$ given by

$$
\begin{equation*}
\sigma_{I}=\gamma\left[1-(3-2 / \gamma)^{2}\right]^{1 / 2} \tag{4}
\end{equation*}
$$

and that a neutral curve (which divides regions of growing and decaying solutions) in the $\gamma-\delta$ plane has

$$
\begin{equation*}
\delta=\arccos (3-2 / \gamma) /\left[\gamma^{2}-(3 \gamma-2)^{2}\right]^{1 / 2} . \tag{5}
\end{equation*}
$$

Deduce that on this curve $\delta \rightarrow \infty$ as $\gamma \rightarrow 1 / 2$ from above.
(c) By considering $\gamma=1-\epsilon$ where $0<\epsilon \ll 1$, or otherwise, show that one possible solution of (5) has $\gamma \delta \rightarrow 1$ as $\gamma \rightarrow 1$. Given that $d \gamma / d \delta$ is negative on this curve, sketch the shape of the curve.
(d) By considering the behaviour of (3) for $\delta=0$, or otherwise, discuss qualitatively the behaviour of solutions near the neutral curve.
Discuss qualitatively the expected behaviour of solutions of (2).
What is the relevance of this behaviour to the ENSO cycle?
4. Two immiscible viscous fluids (say fluid 1 and fluid 2) are in planar motion, touching each other at their common interface, the fluid densities being $\rho_{1}, \rho_{2}$ and viscosities $\mu_{1}, \mu_{2}$ respectively.
(a) Show that the governing equations in fluid 1 may be written in non-dimensional form as

$$
\operatorname{div} \mathbf{u}=0, \quad \partial \mathbf{u} / \partial t+(\mathbf{u} \cdot \operatorname{grad}) \mathbf{u}=-\operatorname{grad} p+\operatorname{Re}^{-1} \Delta \mathbf{u}
$$

where $\mathbf{u}=(u, v)$ is the velocity vector and $p$ is the pressure. Define the Reynolds number Re.
(b) Give the corresponding equations in fluid 2 subject to the same non-dimensionalisation as in (a). Also give the conditions at the interface.
(c) If only a small thin layer of fluid 2 is present, with aspect ratio $h \ll 1$, provide a justification for the expansions

$$
(u, v, p)=\left(u_{1}, v_{1}, h^{-1} p_{1}\right)+\ldots \text { and }(u, v, p)=\left(h^{-1} u_{2}, v_{2}, h^{-1} p_{2}\right)+\ldots
$$

in fluids 1,2 respectively. Provide also expressions for the length and time scales.
(d) If the ratios $\rho_{2} / \rho_{1}$ and $\mu_{2} / \mu_{1}$ are small and comparable, derive from (c) the different systems of equations that can arise within fluid 1 and within the thin layer of fluid 2 over different Re ranges. Identify a critical Re range.
5. A one-dimensional model of viscous blood flow in an artery is governed by the following system of equations

$$
\begin{aligned}
\frac{\partial A}{\partial t}+\frac{\partial Q}{\partial z} & =0 \\
\frac{\partial Q}{\partial t}+\frac{\partial}{\partial z}\left(\frac{Q^{2}}{A}\right) & =-\frac{A}{\rho} \frac{\partial P}{\partial z}-K \frac{Q}{A} \\
\text { and } P-P_{\mathrm{ext}} & =\beta\left(\frac{\sqrt{A}-\sqrt{A_{0}}}{A_{0}}\right)
\end{aligned}
$$

Here, $Q(z, t), A(z, t)$ and $P(z, t)$ are the volume flux, the cross sectional area of the artery, and the fluid pressure respectively, at a distance $z$ along the artery and at time $t$. All other physical parameters present in the model remain constant.
(a) By eliminating the pressure $P(z, t)$, convert the system into the following form

$$
\begin{equation*}
\frac{\partial \mathbf{u}}{\partial t}+\mathbf{H}(\mathbf{u}) \frac{\partial \mathbf{u}}{\partial z}+\mathbf{B}(\mathbf{u})=0 \tag{1}
\end{equation*}
$$

where

$$
\mathbf{u}=\left[\begin{array}{l}
A \\
Q
\end{array}\right], \quad \mathbf{H}(\mathbf{u})=\left[\begin{array}{cc}
0 & 1 \\
\left(\frac{\beta A^{1 / 2}}{2 \rho A_{0}}-\frac{Q^{2}}{A^{2}}\right) & \frac{2 Q}{A}
\end{array}\right], \quad \text { and } \quad \mathbf{B}(\mathbf{u})=\left[\begin{array}{c}
0 \\
K \frac{Q}{A}
\end{array}\right] .
$$

(b) Obtain the eigenvalues $\lambda_{1}$ and $\lambda_{2}$ of the matrix $\mathbf{H}(\mathbf{u})$ and, subsequently, determine a matrix of left eigenvectors $\mathbf{L}$ that satisfies

$$
\mathbf{L} \mathbf{H}(\mathbf{u})=\Lambda \mathbf{L} \quad \text { for } \Lambda=\left[\begin{array}{cc}
\lambda_{1} & 0 \\
0 & \lambda_{2}
\end{array}\right]
$$

(c) By multiplying equation (1) by $\xi(Q, A) \mathbf{L}$, for some smooth scalar function $\xi$, derive the Riemann invariants

$$
\mathbf{W}=\left[\begin{array}{l}
W_{1} \\
W_{2}
\end{array}\right]=\left[\begin{array}{l}
\frac{Q}{A}+4 \sqrt{\frac{\beta}{2 \rho A_{0}}} A^{1 / 4} \\
\frac{Q}{A}-4 \sqrt{\frac{\beta}{2 \rho A_{0}}} A^{1 / 4}
\end{array}\right]
$$

(d) Hence or otherwise, show that these Riemann invariants satisfy the equations

$$
\left(\frac{\partial}{\partial t}+\left(\frac{Q}{A} \pm \sqrt{\frac{\beta}{2 \rho A_{0}}} A^{1 / 4}\right) \frac{\partial}{\partial z}\right)\left(\frac{Q}{A} \pm 4 \sqrt{\frac{\beta}{2 \rho A_{0}}} A^{1 / 4}\right)+K \frac{Q}{A^{2}}=0 .
$$

