

# Bondi's $k$ -calculus

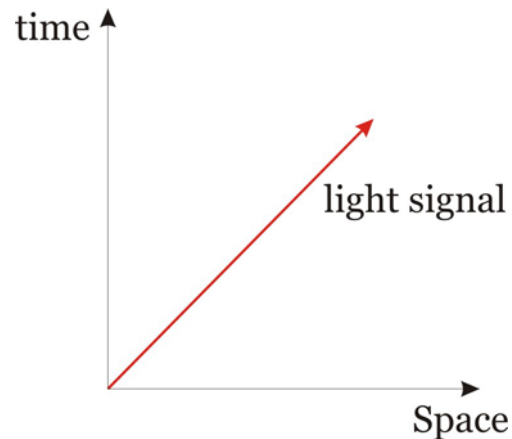
# k-calculus

In order to derive the key results of Special Relativity, we will make use of the k-calculus, invented by Hermann Bondi.



The key idea is to use light signals to relate measurements carried out by observers in (non-accelerated) relative motion.

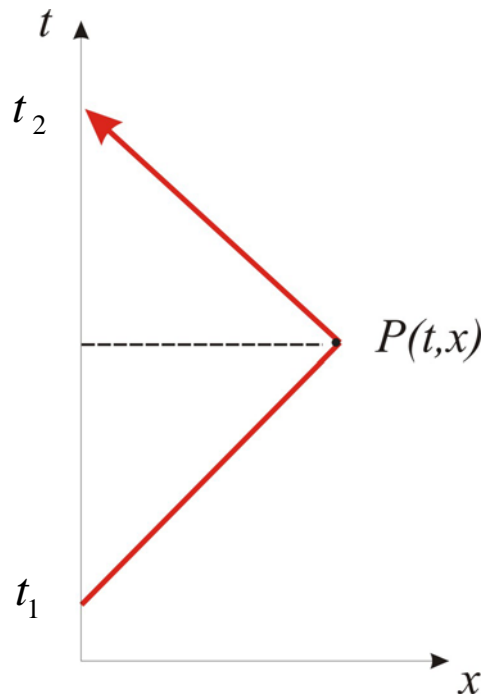
If  $c=1$  then light-signals travel at  $45^\circ$  in a “space-time” diagram.



# Radar method

Idea: Use the constancy of the speed of light to measure distance.

(e.g. time in seconds, distance in light-seconds)



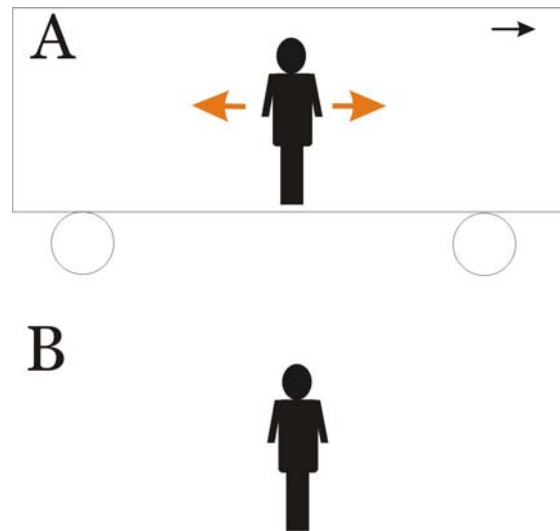
From the diagram it is easy to see that

$$t = \frac{1}{2}(t_1 + t_2)$$

$$x = \frac{1}{2}(t_2 - t_1)$$

# Simultaneity

Consider a simple experiment as viewed by two observers in relative motion.



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According to A:



photons arrive at the same time

According to B:



photon arrives at back before the front

$t=0$

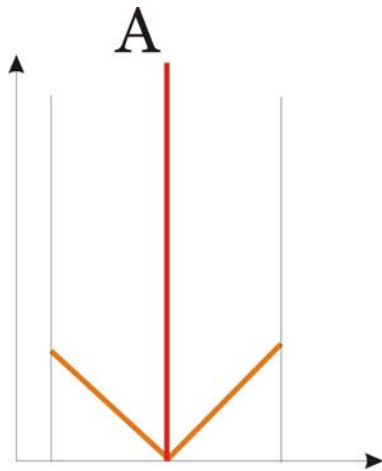
$t=T$

Lesson: Time coordinate of an “event” depends on the observer.

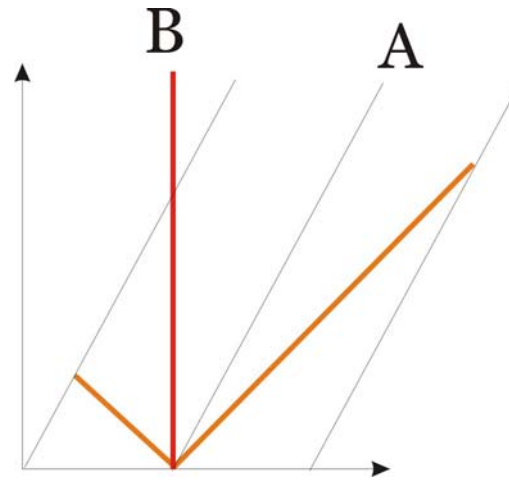
# Simultaneity

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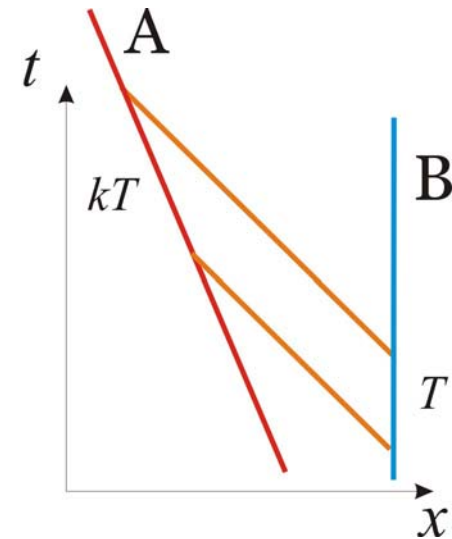
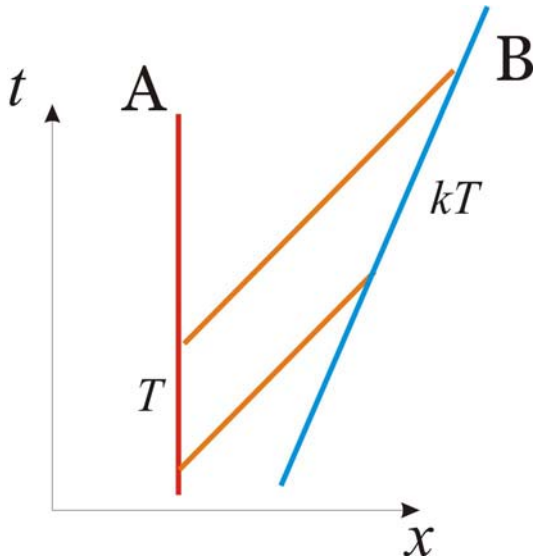
According to A:



According to B:



# The k-factor



Let us now assume that if A observes a time interval  $T$  (say), between two light pulses, then B observes an interval  $kT$ .

Conversely;

If B observes a time interval  $T$  then A observes an interval  $kT$ .

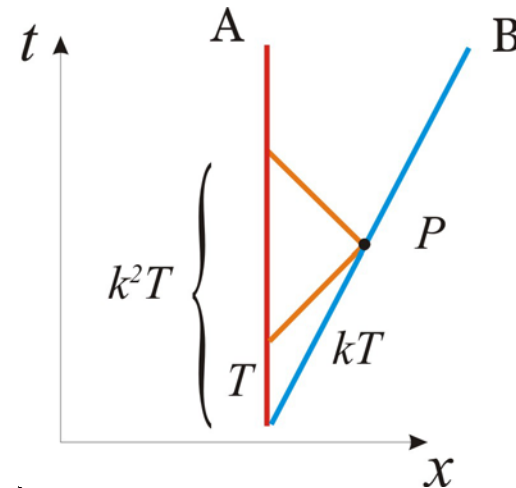
# Relative speed

What are the coordinates of an event  $P$ ?

The “radar method” leads to;

$$t = \frac{1}{2}(T + k^2T) = \frac{1}{2}(k^2 + 1)T$$

$$x = \frac{1}{2}(k^2T - T) = \frac{1}{2}(k^2 - 1)T$$



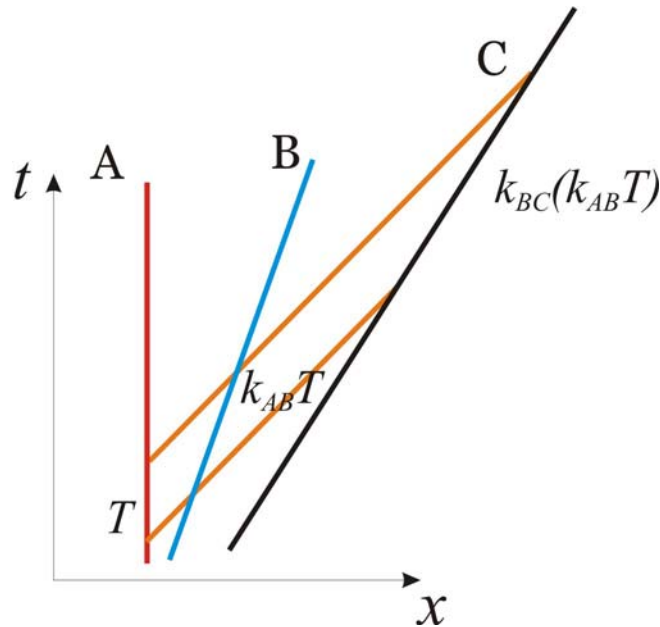
If B has speed  $v$  relative to A, it follows that

$$v = \frac{x}{t} = \frac{k^2 - 1}{k^2 + 1} \quad \Rightarrow \quad k = \left( \frac{1 + v}{1 - v} \right)^{1/2}$$



# Adding velocities

Let us now consider three observer moving relative to one another.



We find that

$$k_{AC} = k_{AB}k_{BC} \quad \Rightarrow \quad v_{AC} = \frac{v_{AB} + v_{BC}}{1 + v_{AB}v_{BC}}$$

That is, velocities do not add in the way that we are used to from Newtonian theory.

“Verified” by Fizeau’s experiment (1851).

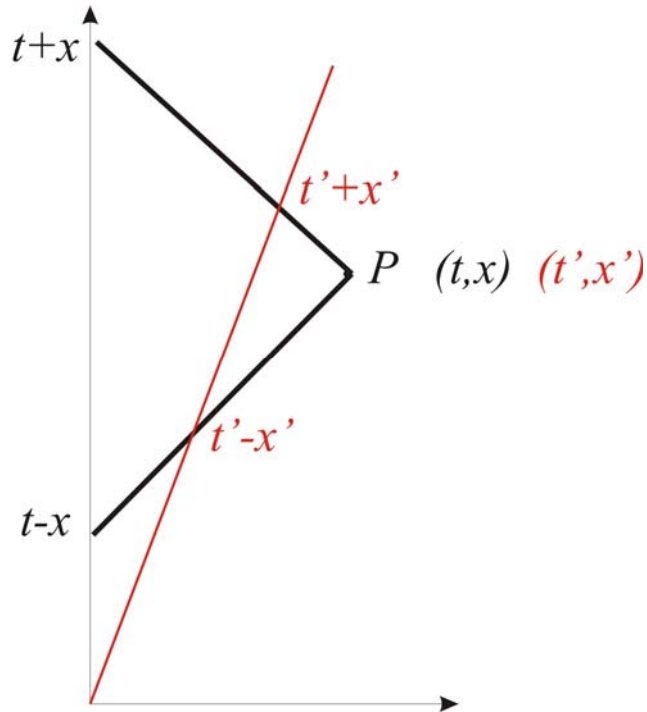
Two limits:

$$v_{AB}, v_{BC} \ll 1 \quad \Rightarrow \quad v_{AC} \approx v_{AB} + v_{BC}$$

$$v_{AB}, v_{BC} \rightarrow 1 \quad \Rightarrow \quad v_{AC} \rightarrow 1$$

# Lorentz transformation

Now consider an event  $P$  viewed by two different observers.



$$t' - x' = k(t - x)$$

$$t + x = k(t' + x')$$

$$t' = \frac{1}{2} \left[ \frac{1}{k} (t + x) + k(t - x) \right] = \frac{t - vx}{(1 - v^2)^{1/2}}$$

$$x' = \frac{1}{2} \left[ \frac{1}{k} (t + x) - k(t - x) \right] = \frac{x - vt}{(1 - v^2)^{1/2}}$$

For low velocities we have

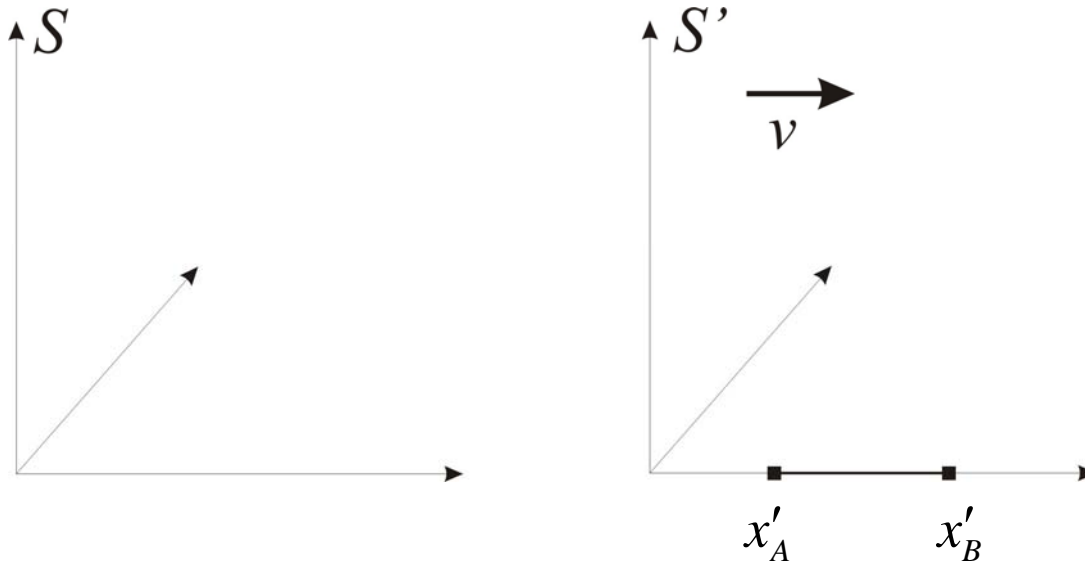
$$t' \approx t \quad \text{and} \quad x' \approx x - vt$$

i.e., a boost in the  $x$ -direction.

# Length contraction

Consider a rod moving with speed  $v$  in the  $x$ -direction.

Let  $S'$  move with the same velocity so that the rod is at rest in this frame.



The rest length (proper length) of the rod is  $l_0 = x'_B - x'_A$

How long does the rod appear to be to observer  $S$ ?

# Length contraction

Measure  $x_A$  and  $x_B$  at the same time  $t_0$ :

$$l = x_B - x_A$$

But we know that

$$x'_A = \beta(x_A - vt_0) \quad \text{where} \quad \beta = (1 - v^2)^{-1/2}$$

$$x'_B = \beta(x_B - vt_0)$$

So we find

$$x'_B - x'_A = \beta(x_B - x_A) \quad \Rightarrow \quad l_0 = \beta l \quad \Rightarrow \quad l < l_0$$

This is an example of length contraction, a moving object appears “shorter” than it is at rest.

# Time dilation

Carry out the same exercise for a clock moving with velocity  $v$ . Let the clock be at rest in the (moving)  $S'$  frame.

Assume that the interval between ticks on the clock is  $T_0$  (proper time). Then ticks correspond to events

$$(x'_1, t'_1) \quad \text{and} \quad (x'_1, t'_1 + T_0)$$

Transforming to the non-moving frame,  $S$ , we have

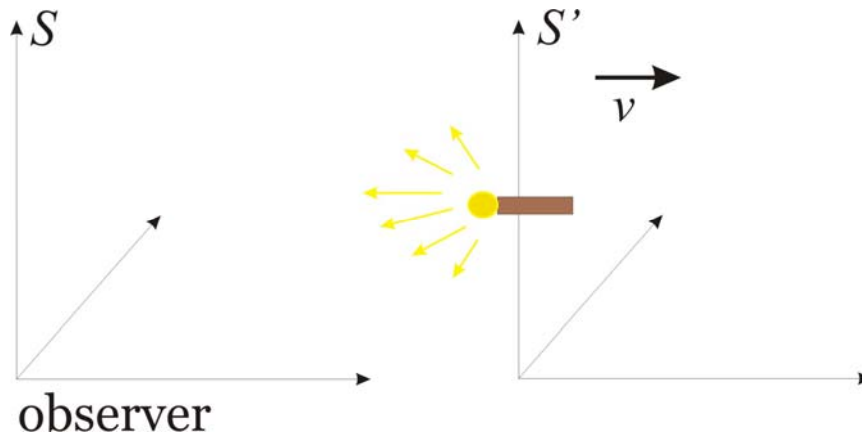
$$\left. \begin{array}{l} t_1 = \beta(t'_1 + vx'_1) \\ t_2 = \beta(t'_1 + T_0 + vx'_1) \end{array} \right\} \Rightarrow T = t_2 - t_1 = \beta T_0 \Rightarrow T > T_0$$

So the moving clock appears to run slower. This is called “time dilation”.

# Doppler effect

Let us first consider the classical (non-relativistic) Doppler effect.

Consider a light source whose wavelength in the rest frame is  $\lambda_0$ . Suppose this source is moving away from an observer with speed  $v$ .



If two photons are emitted a time  $dt'$  apart, then the second photon has to travel an extra distance  $vdt'$ , and hence arrive  $vdt'/c$  later. That is, the photons arrive with a time difference (measured by  $S'$ )

$$\Delta t' = \left(1 + \frac{v}{c}\right) \delta t' \quad \Rightarrow \quad \frac{\lambda}{\lambda_0} = 1 + \frac{v}{c}$$

# Relativistic Doppler shift

Now consider the same problem from the point of view of special relativity.

Due to time dilation, the time difference measured by  $S$  is

$$\delta t = \beta \Delta t' = \beta \left( 1 + \frac{v}{c} \right) \delta t'$$

Since the observed wavelengths are

$$\lambda = c \delta t \quad \text{and} \quad \lambda_0 = c \delta t'$$

we see that

$$\lambda = \beta \left( 1 + \frac{v}{c} \right) \lambda_0 \quad \text{or} \quad \frac{\lambda}{\lambda_0} = \frac{1 + v/c}{\left( 1 - v^2/c^2 \right)^{1/2}}$$

The observed wavelength is longer, the light has been red-shifted.

Note that,

$$\frac{\lambda}{\lambda_0} = \frac{\left( 1 + v/c \right)^{1/2}}{\left( 1 - v/c \right)^{1/2}} = \text{the k-factor}$$

