

Answers for Exam of January 2006

[M283 - Field theory and PDEs]

1. (a) [Standard exercise, similar to lectures and tutorials.]

Gradient:

$$\nabla\phi = \left(\frac{\partial\phi}{\partial x}, \frac{\partial\phi}{\partial y}, \frac{\partial\phi}{\partial z} \right) = (2x, 4y, 2z).$$

2 marks

Directional derivative of Φ :

$$\mathcal{D}_{\mathbf{b}}\Phi = \nabla\Phi \cdot \frac{\mathbf{b}}{|\mathbf{b}|} = (2x, 4y, 2z) \cdot \frac{(1, 0, 0)}{\sqrt{1}}.$$

2 marks

At point $\mathbf{a} = (2, 1, 1)$:

$$\mathcal{D}_{\mathbf{b}}\Phi(\mathbf{a}) = \nabla\Phi(\mathbf{a}) \cdot \frac{\mathbf{b}}{|\mathbf{b}|} = 4.$$

1 marks

Normal to ellipsoid:

$$\hat{\mathbf{n}} = \frac{\nabla\phi}{|\nabla\phi|} = \frac{(x, 2y, z)}{\sqrt{x^2 + 4y^2 + z^2}}.$$

3 marks

At point $\mathbf{a} = (2, 1, 1)$:

$$\hat{\mathbf{n}} = \frac{(2, 2, 1)}{\sqrt{9}}.$$

1 mark

Cartesian equation of tangent plane:

$$\begin{aligned}(\mathbf{r} - \mathbf{a}) \cdot \hat{\mathbf{n}} &= 0 \\ 2(x - 2) + 2(y - 1) + (z - 1) &= 0\end{aligned}$$

3 marks

Or

$$2x + 2y + z = 7.$$

1 mark

- (b) [Bookwork i.e similar to lectures, tutorials and homeworks].

From definition of curl

$$\nabla \times \left(\frac{\mathbf{v}}{\phi} \right) = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \times (\phi^{-1}v_x, \phi^{-1}v_y, \phi^{-1}v_z) \quad (1)$$

1 mark

So that,

$$\begin{aligned}\nabla \times \left(\frac{\mathbf{v}}{\phi} \right) &= \phi^{-2} \left\{ \phi \left[\left(\frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z} \right) \hat{\mathbf{i}} - \left(\frac{\partial v_z}{\partial x} - \frac{\partial v_x}{\partial z} \right) \hat{\mathbf{j}} \right. \right. \\ &+ \left. \left(\frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y} \right) \hat{\mathbf{k}} \right] + \left[\left(v_y \frac{\partial \phi}{\partial z} - v_z \frac{\partial \phi}{\partial y} \right) \hat{\mathbf{i}} \right. \\ &- \left. \left(v_x \frac{\partial \phi}{\partial z} - v_z \frac{\partial \phi}{\partial x} \right) \hat{\mathbf{j}} + \left. \left(v_x \frac{\partial \phi}{\partial y} - v_y \frac{\partial \phi}{\partial x} \right) \hat{\mathbf{k}} \right] \right\} \\ &= \frac{\phi \nabla \times \mathbf{v} + \mathbf{v} \times \nabla \phi}{\phi^2} = \frac{\phi \nabla \times \mathbf{v} - \nabla \phi \times \mathbf{v}}{\phi^2} \quad (2)\end{aligned}$$

3 marks

[This is a standard exercise on gradient and curl discussed in lectures (see chapter 3 of lecture notes) and tutorials 2 and 3.] We check that

$$\begin{aligned}\nabla \left(\frac{1}{r} \right) &= \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) (x^2 + y^2 + z^2)^{-1/2} \\ &= -\frac{1}{r^3} \mathbf{r}, \quad r \neq 0,\end{aligned}$$

2 marks

and

$$\begin{aligned}\nabla(r^3) &= \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) (x^2 + y^2 + z^2)^{3/2} \\ &= 3r\mathbf{r},\end{aligned}$$

1 mark

as well as

$$\begin{aligned}\nabla \times (\mathbf{r}) &= \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \times (x, y, z) \\ &= \mathbf{0},\end{aligned}$$

1 mark

where $\mathbf{r} = x\hat{\mathbf{i}} + y\hat{\mathbf{j}} + z\hat{\mathbf{k}}$ and $r = |\mathbf{r}|$.

[Standard exercise] Therefore, we end up with

$$\begin{aligned}\nabla \times \nabla \left(\frac{1}{r} \right) &= \nabla \times \left(-\frac{\mathbf{r}}{r^3} \right) \\ &= -\frac{r^3(\mathbf{0}) - 3r\mathbf{r} \times \mathbf{r}}{r^6} = \mathbf{0}, \quad r \neq 0.\end{aligned}$$

3 marks

In fact, this result holds for any smooth scalar field ϕ as seen during the lectures [see also exercise 2]. **1 mark**

2. (a) [Standard exercise, similar to lectures and tutorials.]

Line integral of \mathbf{F} over \mathcal{C}_1 :

$$\begin{aligned} \int_{\mathcal{C}_1} \mathbf{F} \cdot d\mathbf{r} &= \int_0^{2\pi} \mathbf{F} \cdot \frac{d\mathbf{r}}{dt} dt \\ &= \int_0^{2\pi} (2 \cos t + t, \cos t + 2t \sin^2 t, 2t^2 \sin t) \cdot (-\sin t, 1, \cos t) dt \\ &= \int_0^{2\pi} t^2 \sin(2t) dt + \int_0^{2\pi} t dt - \int_0^{2\pi} t \sin(t) dt \\ &= 2\pi . \end{aligned}$$

7 marks

- (b) Using

$$\begin{aligned} \nabla \times \nabla \phi &= \left(\frac{\partial^2}{\partial y \partial z} \phi - \frac{\partial^2}{\partial z \partial y} \phi \right) \hat{\mathbf{i}} \\ &\quad - \left(\frac{\partial^2}{\partial x \partial z} \phi - \frac{\partial^2}{\partial z \partial x} \phi \right) \hat{\mathbf{j}} \\ &\quad + \left(\frac{\partial^2}{\partial x \partial y} \phi - \frac{\partial^2}{\partial y \partial x} \phi \right) \hat{\mathbf{k}} , \end{aligned}$$

3 marks

and

$$\frac{\partial^2}{\partial y \partial z} = \frac{\partial^2}{\partial z \partial y} ,$$

we conclude

$$\nabla \times \nabla \phi = \mathbf{0} .$$

1 mark

[Standard exercise, cf. tutorial 2.]

Curl of \mathbf{F}

$$\nabla \times \mathbf{F} = (4yz - 4yz)\hat{\mathbf{i}} - (0 - 0)\hat{\mathbf{j}} + (1 - 1)\hat{\mathbf{k}} = \mathbf{0} . \quad (3)$$

5 marks

We deduce that \mathbf{F} can be expressed as the gradient of a scalar field ϕ . Hence,

$$\mathbf{F} = (2x + y)\hat{\mathbf{i}} + (x + 2yz^2)\hat{\mathbf{j}} + 2y^2z\hat{\mathbf{k}} = \nabla \phi .$$

1 mark

So that

$$\begin{aligned} \phi &= x^2 + xy + f_1(y, z) \\ \phi &= xy + y^2z^2 + f_2(x, z) \\ \phi &= y^2z^2 + f_3(x, y) , \end{aligned}$$

2 marks

where f_1 , f_2 and f_3 are three arbitrary functions which we need specify.
By inspection, we end up with

$$\phi = x^2 + xy + y^2z^2 + C ,$$

2 marks

where C is an arbitrary constant.

(c) Line integral of \mathbf{F} over \mathcal{C}_2 :

$$\begin{aligned} \int_{\mathcal{C}_2} \mathbf{F} \cdot d\mathbf{r} &= \int_0^{2\pi} \mathbf{F} \cdot \frac{d\mathbf{r}}{dt} dt \\ &= \int_0^{2\pi} (2+t, 1, 0) \cdot (0, 1, 0) dt \\ &= \int_0^{2\pi} dt \\ &= 2\pi . \end{aligned}$$

3 marks

[*Bookwork: Gradient theorem and conservative fields.*]

Since \mathbf{F} derives from a scalar field ϕ , the line integrals in (a) and (c) can be reformulated as

$$\begin{aligned} \int_{\mathcal{C}} \mathbf{F} \cdot d\mathbf{r} &= \int_{\mathcal{C}} \nabla\phi \cdot d\mathbf{r} \\ &= \int_0^{2\pi} \frac{d\phi}{dt} dt \\ &= \phi(1, 2\pi, 0) - \phi(1, 0, 0) \\ &= 2\pi . \end{aligned}$$

1 mark

3. [Bookwork, similar exercises were set in lectures, tutorials and homeworks.]

Let us state Gauss's theorem for a differentiable vector field \mathbf{F} defined over a volume τ with bounding surface S .

Gauss's (or divergence's) theorem :

Given a volume τ which is bounded by a piecewise continuous surface S , and a vector function \mathbf{F} which is continuous and has continuous partial derivatives in τ , then

$$\int_{\tau} \nabla \cdot \mathbf{F} d\tau = \oint_S \mathbf{F} \cdot d\mathbf{S} . \quad (4)$$

Notice that the surface integral S has been drawn with a circle around it, in order to indicate that this surface is *closed*, i.e. that it entirely encompasses the volume τ . Also, $d\mathbf{S}$ is defined as the product of a small area (let's say $dx dy$) by the unit outward normal $\hat{\mathbf{n}}$ to the surface S .

5 marks

The exercise now splits in two sections (which can be independently addressed).

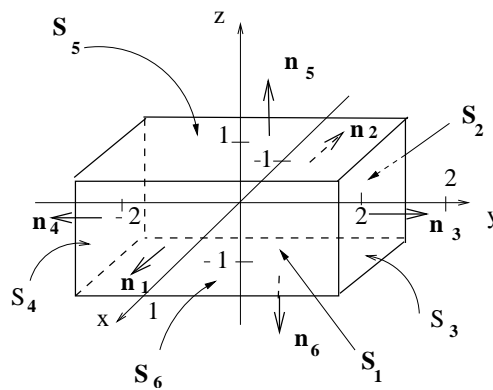


Figure 1: A rectangular box over which we apply Gauss's theorem.

(a) Divergence of \mathbf{F} :

$$\nabla \cdot \mathbf{F} = 1 + 4y + 3 .$$

3 marks

Volume integral of divergence of \mathbf{F} :

$$\int_{\tau} (1 + 4y + 3) d\tau = \int_{-1}^1 dx \int_{-2}^2 dy \int_{-1}^1 dz (4 + 4y) = 32 \int_{-1}^1 dx = 64 . \quad (5)$$

3 marks

By Gauss's theorem, we deduce that

$$\oint_S \mathbf{F} \cdot d\mathbf{S} = \int_{\tau} (\nabla \cdot \mathbf{F}) d\tau = 64 ,$$

where τ is the interior of S .

Let us now check that this is indeed what we would get from the direct evaluation of the surface integral of \mathbf{F} over the surface \mathbf{S} . For this, we decompose \mathbf{S} into six surfaces $\mathbf{S}_1, \dots, \mathbf{S}_6$ oriented by unit normals $\mathbf{n}_1, \dots, \mathbf{n}_6$ as defined in Figure 1.

Let us first evaluate:

$$\begin{aligned}\int_{S_1} \mathbf{F} \cdot d\mathbf{S} &= \int_{S_1} (x, 2y^2, 3z) \cdot (1, 0, 0) dS \\ &= \int_{A_{yz}} x \frac{dydz}{(1, 0, 0) \cdot \hat{\mathbf{i}}} \\ &= x \int_{-2}^2 dy \int_{-1}^1 dz \\ &= 8 ,\end{aligned}$$

since $x = 1$ on the surface S_1 . **1 mark**

Similarly, we find

$$\int_{S_2} \mathbf{F} \cdot d\mathbf{S} = -8 , \quad \text{and} \quad \int_{S_3} \mathbf{F} \cdot d\mathbf{S} = 32 . \quad (6)$$

since $x = -1$ on the surface S_2 and $y = 2$ on S_3 . **2 marks**

Also, we have

$$\int_{S_4} \mathbf{F} \cdot d\mathbf{S} = 32 , \quad \int_{S_5} \mathbf{F} \cdot d\mathbf{S} = 24 , \quad \text{and} \quad \int_{S_6} \mathbf{F} \cdot d\mathbf{S} = -24$$

since $y = -2$ on the surface S_4 , $z = 1$ on S_5 and $z = -1$ on S_6 . **3 marks**

Altogether, we obtain

$$\begin{aligned}\oint_S \mathbf{F} \cdot d\mathbf{S} &= \sum_{i=1}^6 \int_{S_i} \mathbf{F} \cdot d\mathbf{S} \\ &= 64 .\end{aligned}$$

We have therefore checked Gauss's theorem on this example. **1 mark**

(b) Applying Gauss's theorem, we derive that

$$\begin{aligned}\oint_S \mathbf{r} \cdot d\mathbf{S} &= \int_{\tau} \nabla \cdot \mathbf{r} d\tau \\ &= 3 \int_{\tau} d\tau = 3Vol(\tau) .\end{aligned} \quad (7)$$

7 marks

4. [bookwork (same spirit as exercise 3).]

Let us first state the Stokes' theorem for a differentiable vector field \mathbf{F} defined over a surface S bounded by a closed curve \mathcal{C} .

Stokes' theorem:

Given a surface S which is bounded by a piecewise continuous curve \mathcal{C} , and a vector function \mathbf{F} which is continuous and has continuous partial derivatives on S , then

$$\int \int_S (\nabla \times \mathbf{F}) \cdot d\mathbf{S} = \oint_{\mathcal{C}} \mathbf{F} \cdot d\mathbf{r} \quad (8)$$

5 marks

The exercise now splits into two sections (which can be independently addressed).

(a) Curl of \mathbf{F} :

$$\nabla \times \mathbf{F} = (z^2 + x) \hat{\mathbf{i}} - (0 - 0) \hat{\mathbf{j}} + \hat{\mathbf{k}}(-z - 3) . \quad (9)$$

3 marks

Unit normal $\hat{\mathbf{n}}$ to the surface \mathbf{S} (see Figure 2):

$$\hat{\mathbf{n}} = \frac{\nabla(x^2 + y^2 - 2z)}{|\nabla(x^2 + y^2 - 2z)|} = \frac{(x, y, -1)}{\sqrt{x^2 + y^2 + 1}} .$$

(we deduce the orientation of $\hat{\mathbf{n}}$ from the thumb rule). **1 mark**

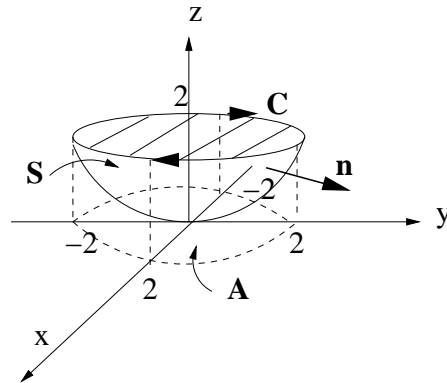


Figure 2: The surface of the paraboloid over which we apply Stokes' theorem.

Thus, an infinitesimal element $d\mathbf{S}$ of orientable surface \mathbf{S} writes as $d\mathbf{S} = \hat{\mathbf{n}} \frac{dxdy}{|\hat{\mathbf{n}} \cdot \hat{\mathbf{k}}|}$ and we end up with

$$\begin{aligned} \int \int_S (\nabla \times \mathbf{F}) \cdot d\mathbf{S} &= \int \int_A (\nabla \times \mathbf{F}) \cdot \hat{\mathbf{n}} \frac{dxdy}{|\hat{\mathbf{n}} \cdot \hat{\mathbf{k}}|} \\ &= \int \int_A \left[x \left(\frac{x^2 + y^2}{2} \right)^2 + x^2 + \frac{x^2 + y^2}{2} + 3 \right] dxdy , \end{aligned}$$

since $z = (x^2 + y^2)/2$ on the surface of the paraboloid.

Note that here, A is the orthogonal projection of \mathbf{S} onto the (xy) -plane as depicted on Figure 2. **3 marks**

Using polar coordinates, we obtain

$$\begin{aligned} \int \int_S (\nabla \times \mathbf{F}) \cdot d\mathbf{S} &= \int_{\theta=0}^{\theta=2\pi} \int_{r=0}^{r=2} [(r \cos \theta)(r^4/2) \\ &+ r^2 \cos^2 \theta + r^2/2 + 3] r dr d\theta = 20\pi . \end{aligned}$$

2 marks

Now, the integral of \mathbf{F} over \mathcal{C} is

$$\oint_{\mathcal{C}} \mathbf{F} \cdot d\mathbf{r} = \int_{2\pi}^0 (-12 \sin^2 t - 8 \cos^2 t) dt = 20\pi . \quad (10)$$

2 marks

Hence, we have checked that

$$\int \int_S (\nabla \times \mathbf{F}) \cdot d\mathbf{S} = 20\pi = \oint_{\mathcal{C}} \mathbf{F} \cdot d\mathbf{r} . \quad (11)$$

(b) Applying Stokes' theorem, we obtain

$$\oint_{\mathcal{C}} \nabla \phi \cdot d\mathbf{r} = \int \int_S (\nabla \times \nabla \phi) \cdot d\mathbf{S} = 0 , \quad (12)$$

since $\nabla \times \nabla \phi = \mathbf{0}$ for any (smooth) scalar function ϕ (see also ex. 2(b)).

7 marks

5. [Standard exercise on Laplace equation i.e. classwork and homework.]

(a) We try a solution which has the form

$$V(x, y) = X(x)Y(y) \quad , \quad (13)$$

which we can rearrange to get

$$\frac{X''(x)}{X(x)} + \frac{Y''(y)}{Y(y)} = 0 \quad . \quad (14)$$

2 marks which leads to (since $Y(0) = Y(L) = 0$)

$$\frac{d^2 X}{dx^2} = +\alpha^2 X(x) \quad , \quad (15)$$

$$\frac{d^2 Y}{dy^2} = -\alpha^2 Y(y) \quad . \quad (16)$$

4 marks

The general solution to the differential equation (16) is

$$Y(y) = C \cos(\alpha y) + D \sin(\alpha y) \quad (17)$$

2 marks

Using $Y(0) = Y(L) = 0$ we are left with

$$\alpha := \alpha_n = n\pi/L \quad \text{where } n = 0, 1, 2, \dots \quad (18)$$

1 mark

We now have the eigenvectors of the differential equation (16):

$$Y_n(y) = D_n \sin(n\pi y/L) \quad , \text{ where } n = 0, 1, 2, \dots \quad (19)$$

1 mark

The general solution to the first differential equation (15) is

$$X(x) = A \cosh(\alpha x) + B \sinh(\alpha x) \quad (20)$$

2 marks

Using $X(0) = 0$, $X(L) = V_0$ we find that the eigenvectors of (15) are

$$X_n(x) = B_n \sinh(\alpha_n x) \quad . \quad (21)$$

2 marks

(b) Combining the functions of x and y , we obtain

$$V_n(x, y) = E_n \sinh(n\pi x/L) \sin(n\pi y/L) , \quad (22)$$

since we already know $\alpha_n = n\pi/L$ from (18).

1 mark

We now use the principle of superposition to write the general solution:

$$V(x, y) = \sum_{n=0}^{\infty} E_n \sinh(n\pi x/L) \sin(n\pi y/L) . \quad (23)$$

1 mark

Now if we pick the coefficients E_n correctly then we can satisfy the final boundary condition, which is

$$V(L, y) = V_0 .$$

So we want to collect coefficients E_n such that

$$V_0 = \sum_{n=1}^{\infty} E_n \sinh(n\pi) \sin(ny) . \quad (24)$$

1 mark

We now multiply equation (24) by the function $\sin(k\pi y/L)$, and integrate it between 0 and L :

$$V_0 \int_0^L \sin(k\pi y/L) dy = E_k \sin(k\pi) \frac{L}{2} , \quad (25)$$

1 mark

where we have used the orthogonality relation (see Hint) to eliminate all terms but one from the infinite sum.

This leads to

$$E_k = \frac{2V_0}{\pi k} \frac{1 - \cos(k\pi)}{\sinh(k\pi)} . \quad (26)$$

1 mark

We notice that $E_{2k} = 0$ and

$$E_{2k-1} = \frac{4V_0}{\pi} \frac{1}{(2k-1) \sinh((2k-1)\pi)} .$$

1 mark

Hence we conclude that

$$\begin{aligned} V(x, y) &= \sum_{n=1}^{\infty} E_n \sinh(n\pi x/L) \sin(n\pi y/L) \\ &= \frac{4V_0}{\pi} \sum_{n=1}^{\infty} \frac{\sinh((2n-1)\frac{\pi y}{L})}{(2n-1) \sinh((2n-1)\pi)} \sin((2n-1)\frac{\pi x}{L}) . \end{aligned} \quad (27)$$

2 marks

(c) [*Aim: To derive the famous Maximum Principle*] At the center of the square, the above expression reduces to

$$\begin{aligned} V(L/2, L/2) &= \frac{4V_0}{\pi} \sum_{n=1}^{\infty} \frac{\sinh((2n-1)\frac{\pi}{2})}{(2n-1) \sinh((2n-1)\pi)} \sin((2n-1)\frac{\pi}{2}) \\ &= \frac{2V_0}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{(2n-1) \cosh((2n-1)\pi/2)}, \end{aligned} \quad (28)$$

where we have used that $\sinh(2s) = 2 \sinh s \cosh s$.

3 marks

Since $\pi/4 = 1 - 1/3 + 1/5 - 1/7 + \dots$ and $\cosh s \geq 1$, it follows that

$$0 \leq V(L/2, L/2) \leq \frac{V_0}{2} \leq V_0.$$

1 mark

Nota Bene: This result holds true everywhere within the square thanks to the so-called maximum principle: The solution of Laplace equation takes its minimum and maximum values on the boundary of the domain.

6. [Standard exercise on the wave equation, classwork and homework.]

(a) The separation of variables leads to

$$\frac{X''(x)}{X(x)} = \frac{1}{c^2} \frac{T''(t)}{T(t)}. \quad (29)$$

2 marks

which by virtue of the boundary conditions $X(0) = X(a) = 0$ gives

$$\frac{d^2 X}{dx^2} = -\alpha^2 X(x) \quad , \quad (30)$$

$$\frac{d^2 T}{dt^2} = -\alpha^2 T(t) \quad . \quad (31)$$

4 marks

The general solution of (30) is:

$$X(x) = A \cos(\alpha x) + B \sin(\alpha x) \quad . \quad (32)$$

1 mark

Using $X(0) = 0$ and $X(a) = 0$ gives us a restraint on α :

$$0 = B \sin(\alpha a) \quad .$$

1 mark

That is,

$$\alpha = \frac{n\pi}{a} \quad \text{where } n = 0, 1, 2, \dots \quad (33)$$

1 mark

We now have the eigenvalues of the differential equation (30):

$$X_n(x) = B_n \sin\left(\frac{n\pi x}{a}\right) \quad , \text{ where } n = 0, 1, 2, \dots \quad (34)$$

1 mark

The general solution to the second differential equation (31) is

$$T(t) = C \cos(\alpha ct) + D \sin(\alpha ct) \quad . \quad (35)$$

1 mark

From (33), we obtain:

$$T_n(t) = C_n \cos\left(\frac{n\pi ct}{a}\right) + D_n \sin\left(\frac{n\pi ct}{a}\right) \quad . \quad (36)$$

1 mark

Combining the functions of x and t and using the principle of superposition we deduce the general solution:

$$V(x, t) = \sum_{n=0}^{\infty} B_n \sin\left(\frac{n\pi x}{a}\right) \left[C_n \cos\left(\frac{n\pi ct}{a}\right) + D_n \sin\left(\frac{n\pi ct}{a}\right) \right]. \quad (37)$$

1 mark

- (b) To fix the constants, we should now use the initial boundary condition $V(x, 0) = 0$ (the displacement of the string is initially zero) which implies that $B_n C_n = 0$. Hence, we deduce that

$$V(x, t) = \sum_{n=1}^{\infty} B_n D_n \sin\left(\frac{n\pi x}{a}\right) \sin\left(\frac{n\pi ct}{a}\right). \quad (38)$$

2 marks

Further, the velocity of the string at $t = 0$ is

$$U = \frac{\partial V}{\partial t}(x, 0) = \sum_{n=1}^{\infty} B_n C_n \sin\left(\frac{n\pi x}{a}\right) \frac{n\pi c}{a}. \quad (39)$$

1 mark

Multiplying equation (39) by the function $\sin\frac{k\pi x}{a}$, and integrating it between 0 and a we obtain:

$$U \int_0^a \sin\frac{k\pi x}{a} dx = \sum_{n=1}^{\infty} B_n C_n \left(\int_0^a \sin\frac{n\pi x}{a} \sin\frac{k\pi x}{a} dx \right) \frac{n\pi c}{a} = B_k C_k \frac{k\pi c}{2} \quad (40)$$

where we have used the orthogonality relation (see Hint) to eliminate all terms but one from the infinite sum.

2 marks

This leads to

$$B_k C_k = -2aU \frac{\cos(k\pi) - 1}{(k\pi)^2 c}. \quad (41)$$

1 mark

We now notice that $B_{2k} C_{2k} = 0$ and

$$B_{2k-1} C_{2k-1} = \frac{4aU}{\pi^2 c} \frac{1}{(2k-1)^2}.$$

We can now write down the full solution to the differential equation:

$$V(x, t) = \frac{4aU}{\pi^2 c} \sum_{n=1}^{\infty} \frac{\sin\frac{(2n-1)\pi x}{a} \sin\frac{(2n-1)\pi ct}{a}}{(2n-1)^2}. \quad (42)$$

1 mark

Using the trigonometric identity

$$\sin(pt) \sin(rt) = 1/2 (\cos((p-r)t) - \cos((p+r)t)) ,$$

we obtain

$$V(x, t) = \frac{2aU}{\pi^2 c} \left[\sum_{n=1}^{\infty} \frac{1}{(2n-1)^2} \cos \left((2n-1)(x-ct) \frac{\pi}{a} \right) - \sum_{n=1}^{\infty} \frac{1}{(2n-1)^2} \cos \left((2n-1)(x+ct) \frac{\pi}{a} \right) \right] . \quad (43)$$

The physical interpretation is that the left series represents wave propagation to the right with speed c and the right series represents wave propagation to the left with speed c .

1 mark