

MATH262001

This question paper consists of 5 printed pages, each of which is identified by the reference **MATH2620**.

Only approved basic scientific calculators may be used.

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Examination for the Module MATH2620

(May 2007)

Fluid Dynamics

Time allowed: **2 hours**

Answer **FOUR** of the **FIVE** questions.

All questions carry equal marks.

1. (a) Explain what is meant by a particle path and a streamline. Under what circumstances are these the same?
 (b) A two-dimensional flow is given by the velocity field

$$\mathbf{u} = \left(\frac{y}{b^2}, -\frac{(x - x_0)}{a^2} \right),$$

where x_0 , a and b are positive constants.

Find the particle paths $(x(t), y(t))$ for this flow for the particle at $(2x_0, 0)$ at $t = 0$.

Show that this fluid flow is incompressible and calculate the corresponding streamfunction $\psi(x, t)$. Hence sketch the streamlines for this flow. Verify that for this flow the streamline through the point $(2x_0, 0)$ is the same as the particle path.

Write down the formula for the acceleration of a fluid particle, and hence calculate the fluid acceleration at a general point (x, y) at time t .

2. State the conditions under which the fluid velocity may be written as the gradient of a velocity potential, ϕ . Show that, when ϕ exists and the flow is incompressible, ϕ satisfies Laplace's equation.

An incompressible, two-dimensional, irrotational flow occupies the half-space $y < 0$. The potential satisfies Laplace's equation and the boundary conditions $\phi = 0$ at $x = 0, 2\pi$, and $\phi = \sin kx$ ($k > 0$) on $y = 0$, and $\phi \rightarrow 0$ as $y \rightarrow -\infty$. Using the method of separable solutions show that

$$\phi = \exp(ky) \sin(kx).$$

Calculate the velocity $\mathbf{u}(x, y)$ and the streamfunction $\psi(x, y)$. Sketch the streamlines for $k = 1$. Show that $(\nabla\phi) \cdot (\nabla\psi) = 0$.

3. Write down the equation of conservation of mass in Cartesian coordinates for an incompressible two-dimensional flow. State the relationship between the streamfunction $\psi(x, y)$ and the fluid velocity, and show that this satisfies the equation of conservation of mass. Show that ψ is constant along a streamline.

Define the vorticity of a velocity field, and show that for a two-dimensional flow

$$\nabla^2\psi = -\omega$$

where the vorticity $\omega = (0, 0, \omega)$.

Two-dimensional flow is set-up between the two coaxial circular cylinders $r = a$ and $r = b$ ($b > a$). The vorticity distribution is given by

$$\omega = \frac{a^3}{4r} - r^2.$$

The azimuthal flow $u_\theta = 0$ on $r = a$ and $u_\theta = \Omega$ on $r = b$. On the assumption that the streamfunction is solely a function of radial distance, r , find the streamfunction and hence the velocity field in polar coordinates. Show that

$$\Omega = \frac{a^3 - b^3}{4}.$$

Calculate directly the circulation

$$\Gamma = \int \mathbf{u} \cdot d\mathbf{x}$$

around the circle $r = b$, and show that it is equal to

$$\int_S \omega dS,$$

where S is the area between $r = a$ and $r = b$.

4. Starting from Euler's equation in the form

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\frac{1}{\rho} \nabla p + \mathbf{g}$$

show that for a steady flow

$$\frac{p}{\rho} + gz + \frac{1}{2}u^2$$

is constant along a streamline, where z is the upward vertical coordinate. (You may quote any of the vector identities at the end of the examination paper).

Show, in addition, that for a steady irrotational flow ($\omega = 0$) that

$$\frac{p}{\rho} + gz + \frac{1}{2}u^2$$

is constant everywhere in the flow.

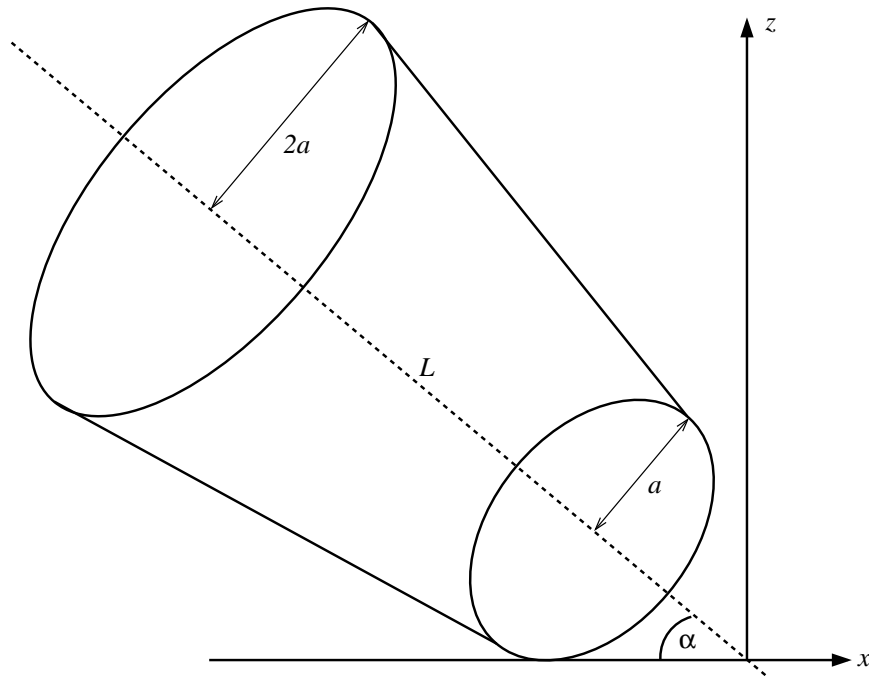


Figure 1: Section of the pipe for Question 4

A long straight pipe of length L has a slowly varying circular cross-section. It is inclined so that its axis is at an angle α to the horizontal, with its smaller cross-section downwards (as shown in the figure). The radius of the pipe at its upper end is $2a$ whilst that at the lower end is a . Water (with density ρ) is pumped at a steady rate through the pipe such that the pressure at the top of the pipe is $2p_a$ and that at the bottom is atmospheric pressure p_a . Show that the water emerges from the pipe with speed U given by

$$U^2 = \frac{32}{15} \left(gL \sin \alpha + \frac{p_a}{\rho} \right),$$

and find the velocity half-way along the pipe.

5. Water flows in a long horizontal channel. The breadth, b , of the channel is equal to B except for an intermediate section where the channel narrows gradually to a minimum width B_m before gradually widening again to B . Upstream the fluid velocity is U and the water is of depth H .

Write down relations obtained from the conservation of mass and Bernoulli's equation.

Define the Froude number, and show that if the upstream Froude number is $\sqrt{3/5}$ then the breadth of the channel, b , and the height of fluid, h , are related by

$$\frac{B^2}{b^2} = \frac{h^2}{H^2} \left(\frac{13}{3} - \frac{10h}{3H} \right)$$

Show that the height of the flow downstream of the constriction must be either $\frac{h}{H} = 1$ or $\frac{h}{H} = \frac{3+\sqrt{129}}{20}$. In the latter case calculate the depth and velocity of the water in the channel at the narrowest point.

Formulae Sheet

Useful Vector Identities

$$\begin{aligned}
 \nabla \times \nabla p &= 0, \\
 \nabla \cdot (\nabla \times \mathbf{u}) &= 0, \\
 \nabla \cdot (p\mathbf{u}) &= p\nabla \cdot \mathbf{u} + \mathbf{u} \cdot \nabla p, \\
 \nabla \times (p\mathbf{u}) &= p\nabla \times \mathbf{u} + \nabla p \times \mathbf{u}, \\
 \nabla \times (\mathbf{B} \times \mathbf{A}) &= \mathbf{A} \cdot \nabla \mathbf{B} - \mathbf{B} \cdot \nabla \mathbf{A} + \mathbf{A} \nabla \cdot \mathbf{B} - \mathbf{B} \nabla \cdot \mathbf{A}, \\
 \nabla \cdot (\mathbf{A} \times \mathbf{B}) &= \mathbf{B} \cdot \nabla \times \mathbf{A} - \mathbf{A} \cdot \nabla \times \mathbf{B}, \\
 \nabla (\mathbf{A} \cdot \mathbf{B}) &= \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A}) + \mathbf{A} \cdot \nabla \mathbf{B} + \mathbf{B} \cdot \nabla \mathbf{A}, \\
 \nabla^2 \mathbf{u} &= \nabla (\nabla \cdot \mathbf{u}) - \nabla \times (\nabla \times \mathbf{u}), \\
 (\nabla \times \mathbf{u}) \times \mathbf{u} &= \mathbf{u} \cdot \nabla \mathbf{u} - \nabla \left(\frac{1}{2} \mathbf{u}^2 \right).
 \end{aligned}$$

Cartesian coordinates

Scalar p , vector $\mathbf{u} = u\mathbf{e}_x + v\mathbf{e}_y + w\mathbf{e}_z$

$$\begin{aligned}
 \text{grad } p &= \nabla p = \frac{\partial p}{\partial x} \mathbf{e}_x + \frac{\partial p}{\partial y} \mathbf{e}_y + \frac{\partial p}{\partial z} \mathbf{e}_z, \\
 \text{div } \mathbf{u} &= \nabla \cdot \mathbf{u} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}, \\
 \text{curl } \mathbf{u} &= \nabla \times \mathbf{u} = \left(\frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right) \mathbf{e}_x + \left(\frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right) \mathbf{e}_y + \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right) \mathbf{e}_z, \\
 \mathbf{u} \cdot \nabla p &= u \frac{\partial p}{\partial x} + v \frac{\partial p}{\partial y} + w \frac{\partial p}{\partial z}, \\
 \nabla^2 p &= \frac{\partial^2 p}{\partial x^2} + \frac{\partial^2 p}{\partial y^2} + \frac{\partial^2 p}{\partial z^2}, \\
 \mathbf{u} \cdot \nabla \mathbf{u} &= \left(u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) \mathbf{e}_x + \left(u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right) \mathbf{e}_y \\
 &\quad + \left(u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right) \mathbf{e}_z
 \end{aligned}$$

Cylindrical Polar Coordinates

$$\mathbf{u} = u\mathbf{e}_r + v\mathbf{e}_\theta + w\mathbf{e}_z$$

$$\begin{aligned}\nabla p &= \frac{\partial p}{\partial r}\mathbf{e}_r + \frac{1}{r}\frac{\partial p}{\partial\theta}\mathbf{e}_\theta + \frac{\partial p}{\partial z}\mathbf{e}_z, \\ \nabla \cdot \mathbf{u} &= \frac{1}{r}\frac{\partial}{\partial r}(ru) + \frac{1}{r}\frac{\partial v}{\partial\theta} + \frac{\partial w}{\partial z}, \\ \nabla \times \mathbf{u} &= \left(\frac{1}{r}\frac{\partial w}{\partial\theta} - \frac{\partial v}{\partial z}\right)\mathbf{e}_r + \left(\frac{\partial u}{\partial z} - \frac{\partial w}{\partial r}\right)\mathbf{e}_\theta + \frac{1}{r}\left(\frac{\partial}{\partial r}(rv) - \frac{\partial u}{\partial\theta}\right)\mathbf{e}_z, \\ \mathbf{u} \cdot \nabla p &= u\frac{\partial p}{\partial r} + \frac{v}{r}\frac{\partial p}{\partial\theta} + w\frac{\partial p}{\partial z}, \\ \nabla^2 p &= \frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial p}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2 p}{\partial\theta^2} + \frac{\partial^2 p}{\partial z^2}.\end{aligned}$$

Spherical Polar Coordinates

$$\mathbf{u} = u\mathbf{e}_r + v\mathbf{e}_\theta + w\mathbf{e}_\phi$$

$$\begin{aligned}\nabla p &= \frac{\partial p}{\partial r}\mathbf{e}_r + \frac{1}{r}\frac{\partial p}{\partial\theta}\mathbf{e}_\theta + \frac{1}{r\sin\theta}\frac{\partial p}{\partial\phi}\mathbf{e}_\phi, \\ \nabla \cdot \mathbf{u} &= \frac{1}{r^2}\frac{\partial}{\partial r}(r^2u) + \frac{1}{r\sin\theta}\frac{\partial}{\partial\theta}(v\sin\theta) + \frac{1}{r\sin\theta}\frac{\partial w}{\partial\phi}, \\ \nabla \times \mathbf{u} &= \frac{1}{r\sin\theta}\left(\frac{\partial}{\partial\theta}(w\sin\theta) - \frac{\partial v}{\partial\phi}\right)\mathbf{e}_r + \frac{1}{r}\left(\frac{1}{\sin\theta}\frac{\partial u}{\partial\phi} - \frac{\partial}{\partial r}(rw)\right)\mathbf{e}_\theta \\ &\quad + \frac{1}{r}\left(\frac{\partial}{\partial r}(rv) - \frac{\partial u}{\partial\theta}\right)\mathbf{e}_\phi, \\ \mathbf{u} \cdot \nabla p &= u\frac{\partial p}{\partial r} + \frac{v}{r}\frac{\partial p}{\partial\theta} + \frac{w}{r\sin\theta}\frac{\partial p}{\partial\phi}, \\ \nabla^2 p &= \frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial p}{\partial r}\right) + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial p}{\partial\theta}\right) + \frac{1}{r^2\sin\theta}\frac{\partial^2 p}{\partial\phi^2}.\end{aligned}$$

Divergence Theorem and Stokes Theorem

Let V be a region bounded by a simple closed surface S with unit **outward** normal \mathbf{n}

$$\int_S \mathbf{u} \cdot \mathbf{n} dS = \int_V \nabla \cdot \mathbf{u} dV, \quad \int_S p \mathbf{n} dS = \int_V \nabla p dV, \quad \int_S \mathbf{u} \times \mathbf{n} dS = - \int_V \nabla \times \mathbf{u} dV.$$

Let C be a simple closed curve spanned by a surface S with unit normal \mathbf{n}

$$\int_C \mathbf{u} \cdot d\mathbf{x} = \int_S (\nabla \times \mathbf{u}) \cdot \mathbf{n} dS, \quad \int_C p d\mathbf{x} = - \int_S (\nabla p) \times \mathbf{n} dS.$$