

# Solutions to 2003 MathII Exam Paper

March 30, 2006

## Section A

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1.1 Consider a change of

$$\vec{r} = x \hat{i} + y \hat{j} + z \hat{k}$$

1

along a  $q_1$ -line ( $q_2, q_3$  - fixed): the partial derivative

2

$$\frac{\partial \vec{r}}{\partial q_1} = \frac{\partial x}{\partial q_1} \hat{i} + \frac{\partial y}{\partial q_1} \hat{j} + \frac{\partial z}{\partial q_1} \hat{k} \equiv h_1 \bar{e}_1$$

will be directed along  $\bar{e}_1$ , where

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$$h_1 = \left| \frac{\partial \vec{r}}{\partial q_1} \right| = \sqrt{\left( \frac{\partial x}{\partial q_1} \right)^2 + \left( \frac{\partial y}{\partial q_1} \right)^2 + \left( \frac{\partial z}{\partial q_1} \right)^2}$$

is the scale factor, introduced to make  $\bar{e}_1$  of length one.

Similarly for  $\bar{e}_2, h_2$  and  $\bar{e}_3, h_3$ :

$$h_i = \left| \frac{\partial \vec{r}}{\partial q_i} \right| \text{ and } \bar{e}_i = \frac{1}{h_i} \frac{\partial \vec{r}}{\partial q_i}$$

1.2

$$x = r \sin \theta \cos \phi, \quad y = r \sin \theta \sin \phi, \quad z = r \cos \theta$$

Using general formulae from question 1.1:

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$$\textcircled{a} \quad \frac{\partial x}{\partial r} = \sin \theta \cos \phi, \quad \frac{\partial y}{\partial r} = \sin \theta \sin \phi, \quad \frac{\partial z}{\partial r} = \cos \theta$$

$$\frac{\partial x}{\partial \theta} = r \cos \theta \cos \phi, \quad \frac{\partial y}{\partial \theta} = r \cos \theta \sin \phi, \quad \frac{\partial z}{\partial \theta} = -r \sin \theta$$

$$\frac{\partial x}{\partial \phi} = -r \sin \theta \sin \phi, \quad \frac{\partial y}{\partial \phi} = r \sin \theta \cos \phi, \quad \frac{\partial z}{\partial \phi} = 0$$

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(b) Therefore,

$$2 \left[ \begin{aligned} h_r &= \left| \frac{\partial \vec{r}}{\partial r} \right| = \sqrt{(\sin\theta \cos\phi)^2 + (\sin\theta \sin\phi)^2 + (\cos\theta)^2} = 1 \\ h_\theta &= \sqrt{(r \cos\theta \cos\phi)^2 + (r \cos\theta \sin\phi)^2 + (-r \sin\theta)^2} = r \\ h_\phi &= \sqrt{(-r \sin\theta \sin\phi)^2 + (r \sin\theta \cos\phi)^2 + 0^2} = r \sin\theta \end{aligned} \right.$$

(c) and thus

$$2 \left[ \begin{aligned} \vec{e}_r &= \frac{1}{h_r} \left[ \frac{\partial x}{\partial r} \hat{i} + \frac{\partial y}{\partial r} \hat{j} + \frac{\partial z}{\partial r} \hat{k} \right] = \sin\theta (\cos\phi \hat{i} + \sin\phi \hat{j}) + \cos\theta \hat{k} ; \\ \vec{e}_\theta &= \cos\theta (\cos\phi \hat{i} + \sin\phi \hat{j}) - \sin\theta \hat{k} \\ \vec{e}_\phi &= -\sin\phi \hat{i} + \cos\phi \hat{j} \end{aligned} \right.$$

$$(1.3) \int_{-\infty}^{\infty} \delta\left(\frac{x}{5} - 1\right) f(x) dx = \left| \begin{array}{l} \text{change variable} \\ y = \frac{x}{5} - 1 \end{array} \right|$$

$$7 \left| \int_{-\infty}^{\infty} \delta(y) f(5y+5) dy = 5 \cdot f(5) \right.$$

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1.4) formal definition of the integral Fourier transform

$$2 \quad F(\nu) = \mathcal{F}[f(t)] = \int_{-\infty}^{\infty} f(t) e^{-i2\pi\nu t} dt$$

~~[+ sign in the exponential is also acceptable]~~

• Inverse transform:

$$2 \quad f(t) = \mathcal{F}^{-1}[F(\nu)] = \int_{-\infty}^{\infty} F(\nu) e^{+i2\pi\nu t} d\nu$$

~~[sign is acceptable if it was chosen in the 1st]~~

[an opposite choice of signs in the exponentials is also perfectly acceptable]

• Consider

$$\mathcal{F}[\delta(t)] = \int_{-\infty}^{\infty} \delta(t) e^{-i2\pi\nu t} dt = 1,$$

3 so that the inverse transform

$$\delta(t) = \int_{-\infty}^{\infty} e^{i2\pi\nu t} d\nu, \text{ G.E.D.}$$

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$$1.5 \quad \underline{1} \quad f(t) = e^{-\alpha|t|} = \begin{cases} e^{-\alpha t}, & t \geq 0 \\ e^{\alpha t}, & t \leq 0 \end{cases}$$

Its Fourier transform

$$F(\nu) = \int_{-\infty}^{\infty} e^{-\alpha|t|} e^{-i2\pi\nu t} dt = \int_{-\infty}^0 e^{\alpha t} e^{-i2\pi\nu t} dt + \int_0^{\infty} e^{-\alpha t} e^{-i2\pi\nu t} dt = \frac{1}{\alpha - i2\pi\nu} e^{(\alpha - i2\pi\nu)t} \Big|_{-\infty}^0 + \frac{1}{-\alpha - i2\pi\nu} e^{-(\alpha + i2\pi\nu)t} \Big|_0^{\infty} =$$

$$6 = \frac{1}{\alpha - i2\pi\nu} + \frac{1}{\alpha + i2\pi\nu} = \frac{2\alpha}{\alpha^2 + (2\pi\nu)^2}$$

1.6 Rewrite equation by dividing on  $(x^2-1)(x-4)$  as:

$$\underline{1} \quad y''(x) + \frac{1}{\underbrace{(x+1)(x-4)}_{p(x)}} y'(x) + \frac{1}{\underbrace{(x-1)(x-4)}_{q(x)}} y(x) = 0$$

There are three singular points:  $q(x)$

(a)  $x = -1$

2  $\left. \begin{array}{l} (x+1)p(x) \text{ is regular} \\ (x+1)^2 q(x) \text{ is regular} \end{array} \right\} \rightarrow \text{regular singular point}$

(b)  $x = +1$

2  $(x-1)p(x), (x-1)^2 q(x)$  are both regular at  $x = 1$   
↳ also is a regular singular point

-0-

②  $x=4$   
again, since both  $(x-4)p(x)$  and  $(x-4)^2 q(x)$  are regular at this point  $\rightarrow$  it is a regular singular point.

1.7) First, checking  $y_1(x+ct) \equiv y_1(u)$ :

$$\frac{\partial y_1}{\partial x} = \frac{dy_1}{du}, \quad \frac{\partial^2 y_1}{\partial x^2} = \frac{d^2 y_1}{du^2} \quad \text{since } \frac{\partial u}{\partial x} = 1;$$

$$3 \quad \frac{\partial y_1}{\partial t} = \frac{dy_1}{du} \frac{\partial u}{\partial t} = \frac{dy_1}{du} \cdot c, \quad \frac{\partial^2 y_1}{\partial t^2} = c^2 \frac{d^2 y_1}{du^2},$$

so that  $\delta y_1 / \delta t^2 = c^2 \cdot \frac{\partial^2 y_1}{\partial x^2}$ , i.e. it is a solution.

Similarly for  $y_2(x,t) = y_2(x-ct) \equiv y_2(v)$ :

$$\frac{\partial y_2}{\partial x} = \frac{dy_2}{dv}, \quad \frac{\partial^2 y_2}{\partial x^2} = \frac{d^2 y_2}{dv^2}$$

$$3 \quad \frac{\partial y_2}{\partial t} = -c \frac{dy_2}{dv}, \quad \frac{\partial^2 y_2}{\partial t^2} = (-c)^2 \frac{d^2 y_2}{dv^2} = c^2 \frac{d^2 y_2}{dv^2} \equiv c^2 \frac{\partial^2 y_2}{\partial x^2}$$

Q.E.D.

1) The sum of two solutions is always a solution.  
(a superposition principle)

1.8 • Definition of the Legendre polynomials. Based on the generating function is as follows:

$$1 \quad G(x,t) = \frac{1}{\sqrt{1-2xt+t^2}} = \sum_{n=0}^{\infty} P_n(x) \cdot t^n$$

• Therefore, what is needed is to expand  $G(x,t)$  in the Maclaurin series with respect to  $t$  for  $h=0$  and  $1$ :

$$2 \quad G(x,t) = \sum_{h=0}^{\infty} \frac{1}{h!} \cdot G^{(h)}(x,t) \Big|_{t=0} \cdot t^h =$$

$$\text{the particular} = G(x,0) + G'(x,0) \cdot t + \dots$$

$$\equiv P_0(x) + P_1(x) \cdot t + \dots$$

so that

$$2 \quad P_0(x) \equiv G(x,0) = 1 ;$$

$$2 \quad P_1(x) \equiv G'(x,0) = \left( \frac{d}{dt} \frac{1}{\sqrt{1-2xt+t^2}} \right) \Big|_{t=0} = -\frac{1}{2} \frac{-2x+2t}{\sqrt{\dots}} \Big|_{t=0} =$$

$$= -\frac{1}{2} \frac{-2x}{1} = x$$

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Section B

② The parabolic coordinates  $(u, v, \theta)$ :

$$x = uv \cos \theta, \quad y = uv \sin \theta, \quad z = \frac{1}{2}(u^2 - v^2)$$

③ Working out the scale factors:

$$\vec{r} = x \hat{i} + y \hat{j} + z \hat{k}$$

$$\int h_u = \left| \frac{\partial \vec{r}}{\partial u} \right| = \sqrt{\left(\frac{\partial x}{\partial u}\right)^2 + \left(\frac{\partial y}{\partial u}\right)^2 + \left(\frac{\partial z}{\partial u}\right)^2} =$$

$$= \sqrt{(v \cos \theta)^2 + (v \sin \theta)^2 + u^2} = \sqrt{v^2 + u^2}$$

$$\int h_v = \left| \frac{\partial \vec{r}}{\partial v} \right| = \sqrt{(u \cos \theta)^2 + (u \sin \theta)^2 + v^2} = \sqrt{u^2 + v^2}$$

$$\int h_\theta = \left| \frac{\partial \vec{r}}{\partial \theta} \right| = \sqrt{(-uv \sin \theta)^2 + (uv \cos \theta)^2} = uv$$

④ Unit base vectors:

$$\int \bar{e}_u = \frac{1}{h_u} \frac{\partial \vec{r}}{\partial u} = \frac{1}{\sqrt{u^2 + v^2}} \left[ v(\cos \theta \hat{i} + \sin \theta \hat{j}) + u \hat{k} \right],$$

$$\int \bar{e}_v = \frac{1}{h_v} \frac{\partial \vec{r}}{\partial v} = \frac{1}{\sqrt{u^2 + v^2}} \left[ u(\cos \theta \hat{i} + \sin \theta \hat{j}) - v \hat{k} \right],$$

$$\int \bar{e}_\theta = \frac{1}{h_\theta} \frac{\partial \vec{r}}{\partial \theta} = \frac{1}{uv} \left[ uv(-\sin \theta \hat{i} + \cos \theta \hat{j}) \right] =$$

$$= -\sin \theta \hat{i} + \cos \theta \hat{j}$$



ⓐ Checking ~~for~~ the orthogonality of  $\bar{e}_u, \bar{e}_v, \bar{e}_\theta$ :

$$\bar{e}_u \cdot \bar{e}_v = \frac{1}{u^2+v^2} [uv(\cos^2\theta + \sin^2\theta) - uv] = 0,$$

$$\bar{e}_u \cdot \bar{e}_\theta = \frac{1}{\sqrt{u^2+v^2}} \cdot \frac{1}{uv} [uv^2(-\cos\theta \sin\theta + \sin\theta \cos\theta)] = 0,$$

$$\bar{e}_v \cdot \bar{e}_\theta = \frac{1}{\sqrt{u^2+v^2}} \cdot \frac{1}{uv} [u^2 \cdot v(-\cos\theta \sin\theta + \sin\theta \cos\theta)] = 0$$

1 It is orthogonal indeed.

ⓓ We should simply apply the result ~~of~~ in the case of the parabolic coordinates:   
 given

$$\left\{ \begin{array}{l} \text{and } q_1 \rightarrow u, \quad q_2 \rightarrow v, \quad q_3 \rightarrow \theta \\ h_1 = h_2 = \sqrt{u^2+v^2}, \quad h_3 = uv \\ \text{so that we obtain:} \end{array} \right.$$

$$\nabla^2 \psi = \frac{1}{(u^2+v^2)uv} \left[ \frac{\partial}{\partial u} \left( uv \frac{\partial \psi}{\partial u} \right) + \frac{\partial}{\partial v} \left( uv \frac{\partial \psi}{\partial v} \right) + \right.$$

$$\left. + \frac{\partial}{\partial \theta} \left( \frac{u^2+v^2}{uv} \frac{\partial \psi}{\partial \theta} \right) \right] =$$

$$= \frac{1}{uv(u^2+v^2)} \left[ v \frac{\partial}{\partial u} \left( u \frac{\partial \psi}{\partial u} \right) + u \frac{\partial}{\partial v} \left( v \frac{\partial \psi}{\partial v} \right) \right] + \frac{1}{(uv)^2} \frac{\partial^2 \psi}{\partial \theta^2}$$

or, cancelling  $1/uv$ , we obtain the desired result.

③

(a) using the definition of the Fourier transform:

9  
1  
1  
4

$$\begin{aligned} \mathcal{F}[f(t) \cos(2\pi\nu_0 t)] &= \int_{-\infty}^{\infty} f(t) \cos(2\pi\nu_0 t) e^{-i2\pi\nu t} dt = \\ &= \frac{1}{2} \int_{-\infty}^{\infty} f(t) [e^{i2\pi\nu_0 t} + e^{-i2\pi\nu_0 t}] e^{-i2\pi\nu t} dt \\ &= \frac{1}{2} \int_{-\infty}^{\infty} f(t) e^{-i2\pi(\nu-\nu_0)t} dt + \frac{1}{2} \int_{-\infty}^{\infty} f(t) e^{-i2\pi(\nu+\nu_0)t} dt \\ &= \frac{1}{2} [F(\nu-\nu_0) + F(\nu+\nu_0)] \end{aligned}$$

(b) Using integration by parts, we have:

$$\begin{aligned} \mathcal{F}[f'(t)] &= \int_{-\infty}^{\infty} f'(t) e^{-i2\pi\nu t} dt = \\ &= f(t) e^{-i2\pi\nu t} \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} f(t) e^{-i2\pi\nu t} (-i2\pi\nu) dt \\ &= (i2\pi\nu) \int_{-\infty}^{\infty} f(t) e^{-i2\pi\nu t} dt = (i2\pi\nu) F(\nu) \end{aligned}$$

9  
2 since  $f(\pm\infty) = 0$ , ~~which is required by the existence of the Fourier transform.~~

• To calculate  $\mathcal{F}[f''(t)]$  there is no need to go through the same again, since the above

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result is valid for any function. In particular

$$f''(t) = (f'(t))', \text{ so that}$$

$$\hookrightarrow \mathcal{F}[f''(t)] = (i2\pi\nu)^2 \mathcal{F}[f'(t)] = (i2\pi\nu)^2 F(\nu)$$

© Consider

$$F(\nu) = \int f(t) e^{-i2\pi\nu t} dt$$

Then,

$$F'(\nu) = \int f(t) (-i2\pi t) e^{-i2\pi\nu t} dt$$

$$2 \quad F'(0) = (-i2\pi) \int_{-\infty}^{\infty} f(t) t dt = (-2\pi i) \mathcal{M}_1$$

Similarly

$$F''(\nu) = (-2\pi i)^2 \int_{-\infty}^{\infty} f(t) t^2 e^{-i2\pi\nu t} dt$$

$$2 \quad F''(0) = (-2\pi i)^2 \int_{-\infty}^{\infty} f(t) t^2 dt = (-2\pi i)^2 \mathcal{M}_2$$

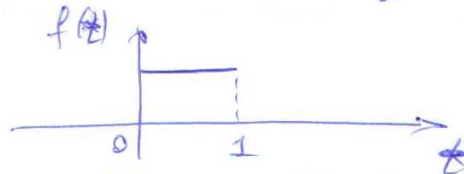
In general, it is easily seen that

$$\hookrightarrow F^{(n)}(\nu) = (-2\pi i)^n \int_{-\infty}^{\infty} f(t) t^n e^{-i2\pi\nu t} dt$$

$$\text{and } F^{(n)}(0) = (-2\pi i)^n \mathcal{M}_n, \text{ Q.E.D.}$$

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$$\textcircled{d} \quad f(x) = \begin{cases} 1, & 0 < x < 1 \\ 0, & \text{for } x < 0 \text{ or } x > 1 \end{cases}$$



$$F[f(x)] = \int_{-\infty}^{\infty} f(t) e^{-i2\pi y t} dt =$$

$$\left. \right\} = \int_0^1 e^{-i2\pi y t} dt = \frac{1}{-i2\pi y} e^{-i2\pi y t} \Big|_0^1 =$$

$$\left. \right\} = \frac{1}{-i2\pi y} (e^{-i2\pi y} - 1) = \frac{1}{i2\pi y} (1 - e^{-i2\pi y})$$

(4)

(a) We rewrite the DE as follows;

1 ~~1~~  $y'' + \frac{1}{x} y' + \left(1 - \frac{p^2}{x^2}\right) y = 0$

1 There is one singular point  $x=0$  which is a regular singular point since ~~for~~

$$\underbrace{x \frac{1}{x}}_{p(x)} = 1 \text{ and } \underbrace{x^2 \left(1 - \frac{p^2}{x^2}\right)}_{q(x)} = x^2 - p^2$$

2 are both regular at  $x=0$ .

(b) To solve the equation, we seek a solution in the

3 form:  $y(x) = x^s \sum_{n=0}^{\infty} a_n x^n$

•  $y' = \sum_{n=0}^{\infty} (n+s) a_n x^{n+s-1}$

2  $y'' = \sum_{n=0}^{\infty} (n+s)(n+s-1) a_n x^{n+s-2}$

•  ~~$y = \sum_{n=0}^{\infty} a_n [(n+s)(n+s-1) x^{n+s} + (n+s) x^{n+s} + \sum_{n=0}^{\infty} a_n x^{n+s+2} - p^2 \sum_{n=0}^{\infty} a_n x^{n+s}] = 0$~~  Substituting into the

DE, we obtain:

$$\sum_{n=0}^{\infty} a_n \left[ (n+s)(n+s-1) x^{n+s} + (n+s) x^{n+s} + \sum_{n=0}^{\infty} a_n x^{n+s+2} - p^2 \sum_{n=0}^{\infty} a_n x^{n+s} \right] = 0$$

3 or  $\sum_{n=0}^{\infty} a_n \left[ (n+s)^2 - p^2 \right] x^{n+s} + \sum_{n=0}^{\infty} a_n x^{n+s+2} = 0$

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- The summation index in the 2nd term is changed:  
 $n \rightarrow n+2$ , while the terms  $n=0$  and  $n=1$  are separated out in the first sum:

$$a_0 (s^2 - p^2) X^s + a_1 [(s+1)^2 - p^2] X^{s+1} + \sum_{n=2}^{\infty} \{ a_n [(n+s)^2 - p^2] + a_{n-2} \} X^{n+s}$$

- The indicial equation reads

$$s^2 - p^2 = 0 \rightarrow s = \pm p$$

there are two roots, corresponding to two solutions.

- Since  $(s+1)^2 - p^2 \neq 0$  for either of the roots, we have to conclude that  $a_1 \equiv 0$  in either case.

- The 3rd term in the eq. above gives the recurrence relation:

$$a_n = - \frac{a_{n-2}}{(n+s)^2 - p^2}$$

- Since  $a_1 = 0$ , it is clear that only even coefficients will survive. ~~The first 3 terms are~~

- To work out the first 3 terms in the solutions, we relate  $a_2$  and  $a_4$  to  $a_0$  using the recurrence relation:

$$a_2 = - \frac{a_0}{(s+2)^2 - p^2} = \begin{cases} - \frac{a_0}{4(p+1)}, & \text{if } s = +p \\ - \frac{a_0}{4(-p+1)}, & \text{if } s = -p \end{cases}$$

$$a_4 = - \frac{a_2}{(s+4)^2 - p^2} = \dots$$

$$2 \left| a_4 = -\frac{a_2}{(4+s)^2 - p^2} = \begin{cases} -\frac{a_2}{8(p+2)} = +\frac{\alpha_0}{8 \cdot 4(p+1)(p+2)}, & \text{if } s = +p \\ -\frac{a_2}{8(-p+2)} = +\frac{\alpha_0}{8 \cdot 4(-p+1)(-p+2)}, & \text{if } s = -p \end{cases} \right.$$

so that the two solutions are:

$$1 \quad y_1(x) = \alpha_0 \left[ 1 - \frac{x^2}{4(p+1)} + \frac{x^4}{8 \cdot 4(p+1)(p+2)} - \dots \right] x^p$$

$$1 \quad y_2(x) = \alpha_0 \left[ 1 - \frac{x^2}{4(-p+1)} + \frac{x^4}{8 \cdot 4(-p+1)(-p+2)} - \dots \right] x^{-p}$$

© The general solution is

$$y(x) = C_1 y_1(x) + C_2 y_2(x)$$

with arbitrary constants  $C_1, C_2$

Note that  $\alpha_0$  is absorbed in both

of them  $\rightarrow$  set to  $\alpha_0 = 1$  in either  
 2 of them.

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⑤ a) We seek a solution of the wave eq.

1 
$$\frac{\partial^2 y}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 y}{\partial t^2}, \quad \underbrace{y(0,t) = y(L,t) = 0}_{\text{the boundary condition}}$$

1 in the form: 
$$y(x,t) = \psi(x) \chi(t)$$

which gives:

$$\psi''(x) \chi(t) = \frac{1}{c^2} \psi(x) \chi''(t)$$

2 
$$\frac{\psi''(x)}{\psi(x)} = \frac{1}{c^2} \frac{\chi''(t)}{\chi(t)} = k \leftarrow \text{the separation constant}$$

thus, we obtain two equations:

2 
$$\underbrace{\psi'' = k\psi}_{\text{involves only } x} \quad \text{and} \quad \underbrace{\chi'' = kc^2\chi}_{\text{involves only } t}$$

⑥ If  $k = -p_n^2 \Rightarrow$  eq. involving  $x$  is

$$\psi'' + p_n^2 \psi = 0$$

2  $\psi_n(x) = \sin p_n x$  satisfies this equation since

$$\psi_n''(x) = -p_n^2 \sin p_n x = -p_n^2 \psi_n(x)$$

It also satisfies the boundary conditions:

2  $\psi_n(x=0) = 0$

$$\psi_n(x=L) = \sin(p_n L) = \sin\left(\frac{2\pi}{L} n \cdot L\right) = \sin 2\pi n = 0$$



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(c) Eq. involving  $t$  is

$$X_n''(t) + c^2 p_n^2 X_n(t) = 0$$

$$X_n(t) = A_n \sin(c p_n t) + B_n \cos(c p_n t)$$

$$X_n'' = -c^2 p_n^2 [A_n \sin(c p_n t) + B_n \cos(c p_n t)]$$

}  $\equiv$

$$-c^2 p_n^2 X_n, \text{ O.E.D.}$$

(d) The general solution is a sum:

$$y(x,t) = \sum_{n=1}^{\infty} \psi_n(x) \otimes X_n(t)$$

$$1 = \sum_{n=1}^{\infty} \sin(p_n x) [A_n \sin(c p_n t) + B_n \cos(c p_n t)]$$

Solutions:

$$y(x,t) = \sum_{n=1}^{\infty} [A_n \sin(c_p n t) + B_n \cos(c_p n t)] \sin(p_n x)$$

with  $A_n, B_n$  - arbitrary constants

2. We use the given initial conditions to obtain the unknown constants  $A_n, B_n$  in the general solution.

• We have:

$$4 \left\{ \begin{aligned} y(x,0) &= \sum_{n=1}^{\infty} B_n \sin(p_n x) \equiv \begin{cases} 0.3 \frac{x}{L}, & 0 \leq x \leq \frac{L}{5} \\ 0.075(1 - \frac{x}{L}), & \frac{L}{5} \leq x \leq L \end{cases} \\ y'_t(x,0) &= \sum_{n=1}^{\infty} c_p n A_n \sin(p_n x) \equiv 0 \end{aligned} \right.$$

2. Multiplying the 2nd equation on  $\sin(p_m x)$  and integrating over  $x$  from 0 to  $L$ , we obtain  $A_m \equiv 0$  since functions  $\sin(p_n x)$  form an orthogonal set.

2. Multiplying the 1st eq. on  $\sin(p_m x)$ , and ~~again~~ integrating over  $x$  and applying the orthogonality relations, we obtain:

$$B_m \frac{L}{2} = \int_0^{\frac{L}{5}} 0.3 \frac{x}{L} \sin p_m x \, dx + \int_{\frac{L}{5}}^L 0.075(1 - \frac{x}{L}) \sin p_m x \, dx$$

Here we use integration by parts:

$$\int_0^{\frac{L}{5}} x \sin p_m x \, dx = \left[ -\frac{1}{p_m} x \cos p_m x + \int \frac{\cos p_m x}{p_m} dx \right]_0^{\frac{L}{5}} =$$

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$$= -\frac{L}{5p_m} \cos \frac{p_m L}{5} + \frac{1}{p_m^2} \sin \frac{p_m L}{5} = -\frac{L^2}{5\pi m} \cos \frac{j\pi m}{5} +$$

$$2 \frac{L^2}{j\pi m^2} \sin \frac{j\pi m}{5},$$

$$\int_{L/5}^L (1 - \frac{x}{L}) \sin p_m x dx = -\left(1 - \frac{x}{L}\right) \frac{\cos p_m x}{p_m} \Big|_{L/5}^L +$$

$$+ \int_{L/5}^L \left(-\frac{1}{L}\right) \frac{\cos p_m x}{p_m} dx = \frac{4}{5p_m} \cos \frac{p_m L}{5} - \frac{1}{L p_m^2} \sin p_m x \Big|_{L/5}^L$$

$$2 = \frac{4L}{5j\pi m} \cos \frac{j\pi m}{5} + \frac{L}{j\pi m^2} \sin \frac{j\pi m}{5}$$

• Collecting all terms, we get:

$$B_m \frac{L}{2} = 0.3 \left[ -\frac{L^2}{5\pi m} \cos \frac{\pi m}{5} + \frac{L^2}{\pi^2 m^2} \sin \frac{\pi m}{5} \right] +$$

$$+ 0.075 \left[ \frac{4L}{5\pi m} \cos \frac{\pi m}{5} + \frac{L}{\pi^2 m^2} \sin \frac{\pi m}{5} \right]$$

$$B_m \frac{L}{2} = 0.375 \frac{L}{\pi^2 m^2} \sin \frac{\pi m}{5}$$

$$1 \quad B_m = \frac{0.75}{\pi^2 m^2} \sin \frac{\pi m}{5}$$

and finally, the partial solution reads:

$$1 \quad y(x,t) = \sum_{n=1}^{\infty} \left[ \frac{0.75}{\pi^2 m^2} \sin \frac{\pi m}{5} \right] \cos\left(\frac{c\pi m}{L} t\right) \cdot \sin\left(\frac{j\pi m}{L} x\right)$$