

**The Handbook of Mathematics, Physics and  
Astronomy Data is provided**

KEELE UNIVERSITY

EXAMINATIONS, 2012/13

Level III

Thursday 2<sup>nd</sup> May 2013, 09.30-11.30

PHYSICS/ASTROPHYSICS

PHY-30028

PHYSICS OF GALAXIES

**Candidates should attempt to answer THREE questions.**

**A sheet of useful formulae can be found on the last page.**

**NOT TO BE REMOVED FROM THE EXAMINATION HALL**

1. (a) Consider a system of main-sequence stars that have individual masses  $M_*$  between the limits  $M_\ell \leq M_* \leq M_u$ , distributed according to  $dN/dM_* = A M_*^{-\beta}$  with  $A$  and  $\beta$  constants.

- i. Assuming the approximate mass–luminosity relation  $L_* \approx M_*^3$  (in Solar units) for individual main-sequence stars, show that the total mass-to-light ratio of the stellar system is

$$\frac{M_{\text{tot}}}{L_{\text{tot}}} \approx \frac{4 - \beta}{2 - \beta} \frac{M_u^{2-\beta} - M_\ell^{2-\beta}}{M_u^{4-\beta} - M_\ell^{4-\beta}} . \quad [20]$$

- ii. Estimate the mass-to-light ratio of an old stellar system with  $\beta = 1.2$ ,  $M_\ell = 0.08 M_\odot$  and  $M_u = 0.9 M_\odot$ . [5]

- (b) Write down the general equation for the acceleration  $\ddot{\mathbf{x}}$  in a gravitational potential  $\Phi(\mathbf{x})$ . Use this to show that the speed on a circular orbit of radius  $R$  in a plane at fixed  $z$  in an axisymmetric galaxy is given by  $V_c^2(R, z) = R (\partial\Phi/\partial R)$ . [20]

- (c) A power-law model for the potential of the Milky Way is

$$\Phi(R, z) = -C (R^2 + z^2/q^2)^{-n} ,$$

where  $C$ ,  $q$  and  $n$  are positive constants.

- i. Show that the circular speed in the midplane ( $z = 0$ ) of the Galaxy, in this model, is  $V_c(R, 0) = \sqrt{2nC} R^{-n}$ . [15]

- ii. At radius  $R$  in an axisymmetric disc, the Oort constants are

$$A \equiv -\frac{1}{2} \left[ \frac{dV_c}{dR} - \frac{V_c}{R} \right] \quad \text{and} \quad B \equiv -\frac{1}{2} \left[ \frac{dV_c}{dR} + \frac{V_c}{R} \right] .$$

Near the Sun, observations give  $A = 14.8 \text{ km s}^{-1} \text{ kpc}^{-1}$  and  $B = -12.4 \text{ km s}^{-1} \text{ kpc}^{-1}$ . Thus, calculate  $V_c$  at the position of the Sun, and infer the value of the exponent  $n$  in a power-law model for the Milky Way potential. [20]

- iii. Within this power-law model, calculate the circular speed at  $R = 50 \text{ kpc}$  in the plane of the Milky Way. Given that the total luminosity of the Galaxy is  $L_{\text{tot}} \approx 2 \times 10^{10} L_\odot$ , estimate the mass of dark matter contained within 50 kpc. [20]

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2. The Plummer model for a self-gravitating, spherical mass distribution has the potential and density

$$\Phi(r) = -\frac{G\mathcal{M}}{(r^2 + a^2)^{1/2}} \quad \text{and} \quad \rho(r) = \frac{3\mathcal{M}}{4\pi} \frac{a^2}{(r^2 + a^2)^{5/2}},$$

where  $\mathcal{M}$  and  $a$  are constants.

- (a) Write down the equation used to obtain  $\rho(r)$  from  $\Phi(r)$  (no calculation is required). [5]

- (b) Show that the mass contained within radius  $r$  in a Plummer sphere is

$$M(r) = \mathcal{M} \frac{r^3}{(r^2 + a^2)^{3/2}}. \quad [15]$$

- (c) Show that the half-mass radius of a Plummer sphere is

$$r_h \simeq 1.3048 a. \quad [15]$$

- (d) The total gravitational potential energy of a spherical mass distribution is given in general by

$$W = \frac{1}{2} \int_0^\infty \Phi(r) \rho(r) 4\pi r^2 dr.$$

- i. Explain this equation in physical terms. [10]

- ii. Show that, for the Plummer model,

$$W \simeq -0.3843 \frac{G\mathcal{M}^2}{r_h}. \quad [25]$$

- (e) The globular star cluster  $\omega$  Centauri has a  $V$ -band luminosity  $L_V = 1.3 \times 10^6 L_\odot$ ; a half-mass radius  $r_h = 9.7$  pc; and an average one-dimensional stellar velocity dispersion  $\sigma = 12$  km s<sup>-1</sup>. The dwarf galaxy Sculptor has  $L_V = 2.3 \times 10^6 L_\odot$  and  $\sigma = 9.9$  km s<sup>-1</sup>, similar to  $\omega$  Centauri, but  $r_h = 380$  pc.

Assuming that the stars in both  $\omega$  Centauri and Sculptor are old, and that Plummer models describe the mass distributions of both systems, use the virial theorem to estimate the mass of dark matter in each. [30]

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3. (a) What is meant by the Eddington limit? Under what conditions might the Eddington luminosity be exceeded? [15]

(b) An accreting black hole of mass  $M$  has a luminosity  $L$ .

i. Write down equations for the forces acting on a proton–electron plasma at a radius of  $r$  from the black hole. [10]

ii. Thus, show that the luminosity ( $L_{\text{Edd}}$ ) of an accreting black hole of mass  $M$  at the Eddington limit is given by

$$L_{\text{Edd}} = \frac{4\pi GcMm_p}{\sigma_T} ,$$

where  $\sigma_T$  is the Thomson cross section. [10]

(c) Consider a particle of mass  $m$  in an accretion disc at a radius  $r$  from a black hole of mass  $M$ .

i. Write down an expression for the gravitational potential energy of the particle. [5]

ii. Use the virial theorem to show that the total luminosity  $L$  emitted by a disc around the black hole is

$$L = \frac{GM\dot{M}}{2r} ,$$

where  $\dot{M}$  is the accretion rate onto the black hole. [15]

(d) The black hole in the quasar 3C 273 has a mass of  $M = 10^9 M_\odot$  and accretes at the Eddington limit.

i. Estimate the luminosity of 3C 273. [5]

ii. Calculate the mass accretion rate of 3C 273 (in Solar units) if the accretion disc extends down to the innermost stable circular orbit (ISCO) of a Schwarzschild black hole. [20]

iii. What is the corresponding efficiency factor of converting rest mass into energy in this case? [10]

iv. How might the efficiency differ for a Kerr black hole? [10]

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4. (a) A virialised, three-dimensional distribution of clouds in the Broad Line Region (BLR) of an AGN has a line-of-sight velocity dispersion  $\sigma$ . If  $r$  is the radial distance of the clouds from the black hole in the AGN, show that the mass of the black hole is

$$M_{\text{BH}} = \frac{3\sigma^2 r}{G} . \quad [15]$$

- (b) The broad  $\text{H}\beta$  line of the Seyfert galaxy, NGC 5548, is measured in a reverberation observation to determine its black hole mass.

- i. The  $\text{H}\beta$  line in NGC 5548 has an emitted wavelength of 486.1 nm and has a Full Width Half Maximum of 8.10 nm. Calculate the velocity dispersion of the BLR clouds. [15]
- ii. The variations in the  $\text{H}\beta$  line are delayed with respect to the AGN continuum. Explain why this delay occurs. [10]
- iii. If the time delay between the continuum and  $\text{H}\beta$  emission is 10 days, estimate the radius of the BLR and hence calculate the mass of the black hole in NGC 5548. [10]

- (c) Suppose that the *stars* around a black hole of mass  $M_{\text{BH}}$  have the density distribution  $\rho_{\star}(r) = K r^{-2}$ , with  $K$  a constant.

- i. Solve the spherical Jeans equation to show that the one-dimensional velocity dispersion of the stars at radius  $r$  from the black hole is given by

$$\sigma_{\star}^2(r) = 2\pi G K + \frac{1}{3} \frac{G M_{\text{BH}}}{r} . \quad [25]$$

- ii. Taking  $K = 1.5 \times 10^5 M_{\odot} \text{ pc}^{-1}$  for stars near the centre of NGC 5548, and using the value of  $M_{\text{BH}}$  from part (b), calculate the radius  $r_{\text{eq}}$  at which the contribution of the black hole to  $\sigma_{\star}$  equals the contribution of the stars themselves. Evaluate  $\sigma_{\star}$  at some appropriate radii in order to make a sketch of  $\sigma_{\star}$  versus  $r$  with scales on both axes. [25]

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5. The hot diffuse plasma in a galaxy cluster can emit X-ray radiation via Bremsstrahlung, whereby the total luminosity  $L$  is given by

$$L = 1.4 \times 10^{-40} Z^2 T^{0.5} \int n_e^2 dV ,$$

where  $Z$  is the average atomic number of the gas,  $T$  is its temperature and  $n_e$  is the electron density. The integral  $\int n_e^2 dV$  is referred to as the *emission measure* of the plasma.

- (a) i. State how Bremsstrahlung radiation is emitted. Why is the emission measure proportional to  $n_e^2$ ? [15]  
 ii. Explain how the luminosity of the X-ray Bremsstrahlung radiation might be affected by cluster mass. [15]
- (b) The electron density of a cluster varies with radius  $r$  as a power-law of the form  $n_e(r) = n_c(r/r_c)^{-\alpha}$ , where  $n_c$  is the core density and  $r_c$  is the core radius of the cluster.

Show that the emission measure of the cluster is equal to

$$\int_0^R n_e^2 dV = \frac{4\pi n_c^2 r_c^{2\alpha}}{(3-2\alpha)} R^{(3-2\alpha)} . \quad [15]$$

- (c) The Coma cluster has a core radius of  $r_c = 200$  kpc,  $kT = 8$  keV and an X-ray luminosity of  $L = 10^{38}$  W, while  $\alpha = 1.3$ .
- i. Evaluate the emission measure of the cluster out to a radius of  $R = 2$  Mpc and hence determine the core electron density (assume the gas is composed of hydrogen). [20]  
 ii. Calculate the total mass of X-ray emitting gas in the cluster inside a radius of 2 Mpc. [25]
- (d) If the total  $B$ -band luminosity of the cluster is  $L_B = 5 \times 10^{12} L_\odot$ , estimate what fraction of the total baryonic mass consists of hot X-ray emitting gas. [10]

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### Some useful formulae

In spherical polar coordinates,

$$\begin{aligned}\ddot{\mathbf{x}} = & \hat{\mathbf{r}} \left( \ddot{r} - r\dot{\theta}^2 - r\sin^2\theta\dot{\phi}^2 \right) + \hat{\boldsymbol{\theta}} \left( r\ddot{\theta} + 2\dot{r}\dot{\theta} - r\sin\theta\cos\theta\dot{\phi}^2 \right) \\ & + \hat{\boldsymbol{\phi}} \left( r\sin\theta\ddot{\phi} + 2\sin\theta\dot{r}\dot{\phi} + 2r\cos\theta\dot{\theta}\dot{\phi} \right)\end{aligned}$$

$$\nabla\Phi = \hat{\mathbf{r}} \frac{\partial\Phi}{\partial r} + \hat{\boldsymbol{\theta}} \frac{1}{r} \frac{\partial\Phi}{\partial\theta} + \hat{\boldsymbol{\phi}} \frac{1}{r\sin\theta} \frac{\partial\Phi}{\partial\phi}$$

$$\nabla^2\Phi = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial\Phi}{\partial r} \right) + \frac{1}{r^2\sin\theta} \frac{\partial}{\partial\theta} \left( \sin\theta \frac{\partial\Phi}{\partial\theta} \right) + \frac{1}{r^2\sin^2\theta} \frac{\partial^2\Phi}{\partial\phi^2}$$

In cylindrical polar coordinates,

$$\ddot{\mathbf{x}} = \hat{\mathbf{R}} \left( \ddot{R} - R\dot{\theta}^2 \right) + \hat{\boldsymbol{\theta}} \left( R\ddot{\theta} + 2\dot{R}\dot{\theta} \right) + \hat{\mathbf{z}} \ddot{z}$$

$$\nabla\Phi = \hat{\mathbf{R}} \frac{\partial\Phi}{\partial R} + \hat{\boldsymbol{\theta}} \frac{1}{R} \frac{\partial\Phi}{\partial\theta} + \hat{\mathbf{z}} \frac{\partial\Phi}{\partial z}$$

$$\nabla^2\Phi = \frac{1}{R} \frac{\partial}{\partial R} \left( R \frac{\partial\Phi}{\partial R} \right) + \frac{1}{R^2} \frac{\partial^2\Phi}{\partial\theta^2} + \frac{\partial^2\Phi}{\partial z^2}$$

Spherical Jeans equation:  $\frac{d}{dr} (\rho\sigma^2) = -\rho(r) \frac{GM(r)}{r^2}$

Standard integrals:  $\int_0^\infty \frac{x^2}{(1+x^2)^3} dx = \frac{\pi}{16}$

$$\int \frac{x^2}{(1+x^2)^{5/2}} dx = \frac{1}{3} \frac{x^3}{(1+x^2)^{3/2}}$$