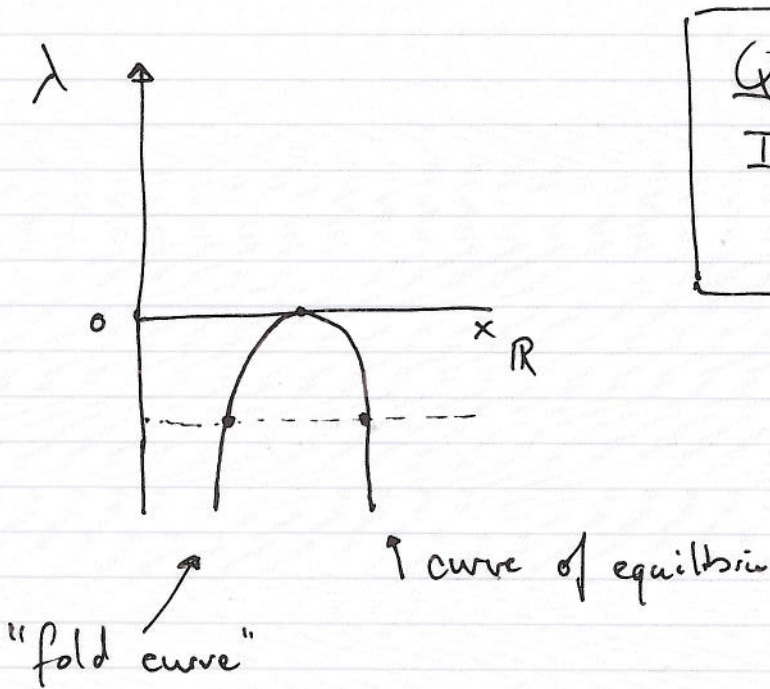
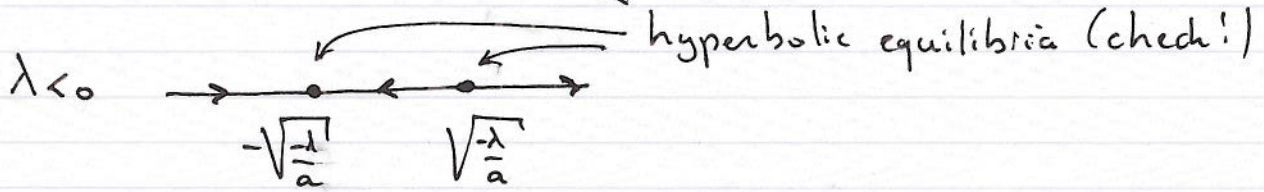
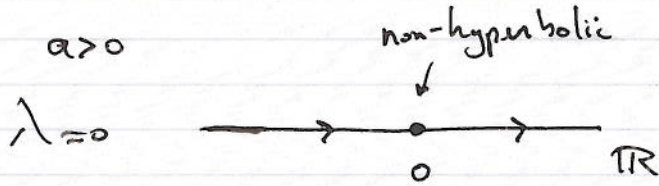


$$\frac{dx}{dt} = f(x, \lambda) \quad x \in \mathbb{R}$$

$$f(x, \lambda) = \lambda + ax^2$$

equilibria $x^2 = -\frac{\lambda}{a}$

if $\lambda < 0$ then no equilibria



Question:
Is this type of bifurcation typical?

Let us consider, more generally

$$f(x, \lambda) = a(\lambda) + b(\lambda)x + c(\lambda)x^2 + O(x^3) \quad x \in \mathbb{R}$$

(general smooth vector field)

Some assumptions: (1) $f(0, 0) = 0$ equilibrium $x=0$ at $\lambda=0$

(2) $D_x f(0, 0) = 0$ equil. is non-hyperbolic

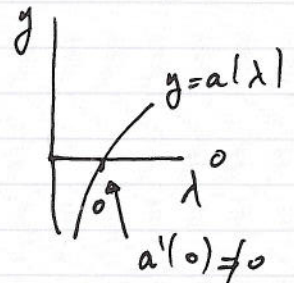
(1) $\Leftrightarrow a(0) = 0$ \uparrow "typical" in 4-par. family of 1D vector fields

(2) $\Leftrightarrow b(0) = 0$

We add another condition (3) $a'(0) \neq 0$ ($\lambda a(\lambda)$ passes 0 with "nonzero" speed)

$$\Rightarrow \lim_{\lambda \rightarrow 0^-} \text{sign } a(\lambda) \neq \lim_{\lambda \rightarrow 0^+} \text{sign } a(\lambda)$$

This is a "typical" property of $a(\lambda)$ as it means that a intersects the line $y=0$ transversely.
 $|\lambda|$ suff small \nearrow the graph of a



Condition (3) implies that

$$D_x^2 f(0, 0) \neq 0$$

$$D_x^2 f(0, 0) = a'(0)$$

with $\lambda(0) = 0$

Then by application of the IFT we have $\exists \lambda(x)$ such that

$$f(x, \lambda(x)) = 0 \quad \text{for } |x| \text{ suff small}$$

Question: what does $\lambda(x)$ look like?

Ans: find the Taylor expansion of λ as fcn of x

(up to certain order)

$$f(x, \lambda) = a(\lambda) + b(\lambda)x + c(\lambda)x^2 + \dots$$

$$\text{w/ } a(\lambda) = a\lambda + o(\lambda^2)$$

→ sloppy $\lambda \approx - \frac{(b(0)x + c(0)x^2)}{a}$ as in simple example.

how to do this more precisely: differentiate $f(x, \lambda) = 0$ wrt x at $x=0$ $\lambda(x)$

$$\frac{d}{dx} f(x, \lambda(x)) \Big|_{x=0} = \frac{d}{dx} (a(\lambda(x)) + b(\lambda(x))x + c(\lambda(x))x^2 + o(x^3)) \Big|_{x=0}$$

$$= a'(0)\lambda'(0) + \cancel{b(0)} = 0 \Rightarrow \boxed{\lambda'(0) = 0}$$

$$\text{now } \frac{d^2}{dx^2} f(x, \lambda(x)) \Big|_{x=0} = \underbrace{a'(0)}_{\neq 0} \lambda''(0) + \cancel{b'(0)\lambda'(0)} + 2c(0) = 0$$

$$\lambda''(0) = - \frac{2c(0)}{a'(0)}$$

if $c(0) \neq 0$ then $\lambda''(0) \neq 0$

⇒ first term a Taylor exp of λ a fn of x is quadratic!

④ if $c(0) \neq 0 \Rightarrow$ "quadratic fold"

if $c(0) = 0$ then maybe $\lambda \approx c \cdot x^3$

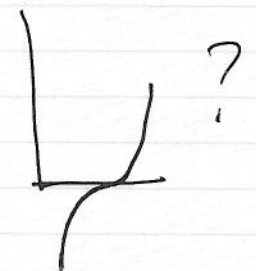
Remember:

* can analyze "typical" systems

* use IFT

* approx soln by Taylor series obtained by diff of formula.

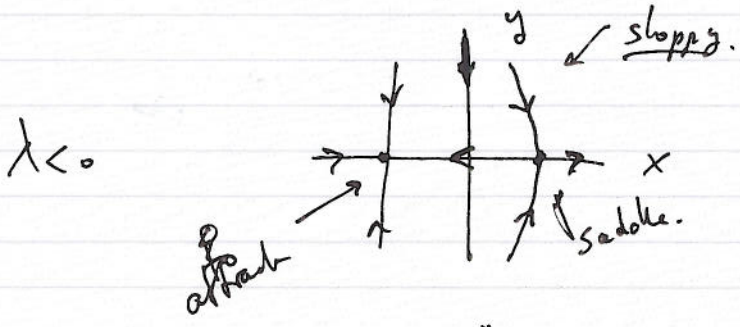
Conclusion: fold is typical bif in one-parameter families of 1-D vector fields.



Does this tell us anything about bifurcations in higher dimensional phase space?

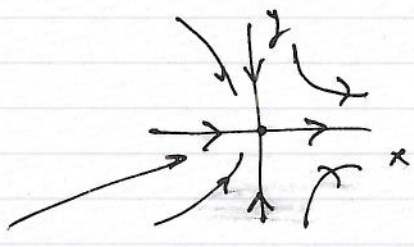
Example:

$$\begin{cases} \dot{x} = \lambda + x^2 \\ \dot{y} = -y \end{cases} \quad \text{ODE on } \mathbb{R}^2$$

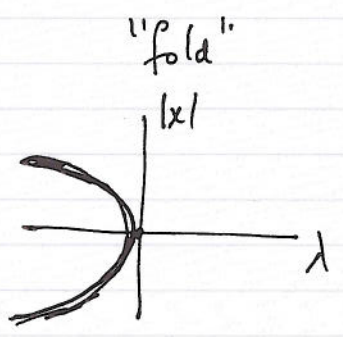


$\lambda = 0$

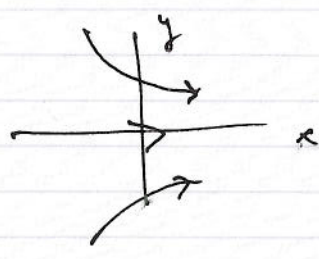
$$Df|_{(0,0)} = \begin{pmatrix} 0 & 0 \\ 0 & -1 \end{pmatrix}$$



but diagram is still



$\lambda > 0$



Observation: the bifurcation takes place in the "non-hyperbolic" direction.

in fact:

Theorem (no proof)

Also if equilibrium is non-hyperbolic, there exist stable and unstable manifolds tangent to (and with the dimension of) E^s and E^u respectively.

NB: there actually also exist a "centre manifold" tangent to E^{so} on which all bifurcations take place. (here x-axis)

Question: in a one-parameter family of vector fields $\dot{x} = f(x, \lambda)$ in \mathbb{R}^m
how can the set of equilibria typically change?

$$\dot{x} = f(x, \lambda) \quad x \in \mathbb{R}^m, \lambda \in \mathbb{R}$$

- Suppose

$$f(0, 0) = 0 \quad x=0 \text{ equilibrium at } \lambda=0.$$

$$\begin{array}{c} \uparrow \uparrow \\ x \quad \lambda \end{array}$$

- if $D_x f(0, 0)$ is invertible $\Rightarrow x=0$ is hyperbolic and the set of equilibria cannot change (locally) if λ is changed because of IFT

$$J^1 x(\lambda) \text{ s.t. } f(x(\lambda), \lambda) = 0 \text{ with } x(0) = 0.$$

- hence we need $D_x f(0, 0)$ to have a zero eigenvalue.

Assume: $D_x f(0, 0)$ has exactly 1 zero eigenvalue.

(this is a persistent phenomenon in 1-par fam. of vector fields)

Question: how to proceed from here?

solving $f(x, \lambda) = 0$ near ~~$x=0, \lambda=0$~~ $x=0, \lambda=0$.

We will try using the IFT (implicit fn thm), by writing $f(x, \lambda) = 0$ as two equations, one that can be solved by IFT and one that cannot be solved this way.

Some observations:

- (1) $\ker D_x f(0, 0)$ has dimension 1 (spanned by eigenvector for eigenvalue 0)

we write:

$$(2) \quad \mathbb{R}^m = \ker D_x f(0, 0) \oplus \mathbb{C}$$

$$\mathbb{C} \cong \mathbb{R}^{m-1}$$

direct sum? \uparrow a complement to $\ker D_x f(0, 0)$ in \mathbb{R}^m

every vector in \mathbb{R}^m can be written uniquely as a sum of one vector in $\ker D_x f(0, 0)$ and one in \mathbb{C}

now introduce new coordinates $\mathbb{R}^m \ni x = (y, z)$
 $\uparrow \quad \uparrow$
 $\in \ker D_1 f(0,0) \quad \in \mathbb{C}$

(3) the Range $D_1 f(0,0)$ has dimension $m-1$

we can write

$$\mathbb{R}^m = \underset{\substack{\uparrow \\ \text{dimension } m-1}}{\text{Range } D_1 f(0,0)} \oplus \underset{\substack{\uparrow \\ \text{complement to Range } D_1 f(0,0) \\ \text{in } \mathbb{R}^m}}{\tilde{\mathbb{C}}}$$

$\tilde{\mathbb{C}} \cong \mathbb{R}$

(4) rewrite ~~f~~ $f: \mathbb{R}^m \times \mathbb{R} \rightarrow \mathbb{R}^m$

$$\text{as } \begin{cases} f_1: \ker D_1 f(0,0) \times \mathbb{C} \times \mathbb{R} \rightarrow \text{Range } D_1 f(0,0) \\ f_2: \ker D_1 f(0,0) \times \mathbb{C} \times \mathbb{R} \rightarrow \tilde{\mathbb{C}} \end{cases}$$

$$f(x, \lambda) = 0 \Leftrightarrow \begin{cases} f_1(y, z, \lambda) = 0 \\ f_2(y, z, \lambda) = 0 \end{cases}$$

(5) The important observation now is that

$D_2 f_1(0,0,0)$ is invertible by construction
 \uparrow
 der. wrt z variable

hence we can use the IFT to solve for z as fn of y, λ
 so that $f_1(y, z(y, \lambda), \lambda) = 0$ with $z(0,0) = 0$

(6) it now remains to solve

$$g(y, \lambda) := f_2(y, z(y, \lambda), \lambda) = 0$$

$$g: \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R} \text{ with } g(0,0) = 0 \quad D_2 g(0,0) = 0$$

We know about this problem:
 (Lyapunov - Schmidt reduction)

$$f(\vec{x}, \lambda) = (\lambda + x^2, -y)^T$$

$$f(0, 0, 0) = 0$$

$$J = D_{x,y} f(0, 0, 0) = \begin{pmatrix} 2x & 0 \\ 0 & -1 \end{pmatrix}$$

$$\mathbb{R}^2 = \ker J \oplus C \quad \ker J = \left\langle \begin{pmatrix} 1 \\ 0 \end{pmatrix} \right\rangle$$

$\begin{matrix} \uparrow & \uparrow \\ \vec{x} & (x, y) \end{matrix}$

$$\text{choose } C = \left\langle \begin{pmatrix} 0 \\ 1 \end{pmatrix} \right\rangle$$

$$\mathbb{R}^2 = \text{Range } J \oplus \tilde{C} \quad \text{Range } J = \left\langle \begin{pmatrix} 0 \\ 1 \end{pmatrix} \right\rangle$$

$$\text{choose } \tilde{C} = \left\langle \begin{pmatrix} 1 \\ 0 \end{pmatrix} \right\rangle$$

$$f_2(x, y, \lambda) = \lambda + x^2 \quad f_2: \ker J \times C \times \mathbb{R} \rightarrow \text{Range } J \oplus \tilde{C}$$

$$\left(D_2 f_1(0, 0, 0) = -1 \right) \quad f_1(x, y, \lambda) = -y \quad f_1: \ker J \times C \times \mathbb{R} \rightarrow \text{Range } J$$

$$\exists! y(x, \lambda) \text{ solving } f(x, y(x, \lambda), \lambda) = 0$$

$$y(x, \lambda) = 0$$

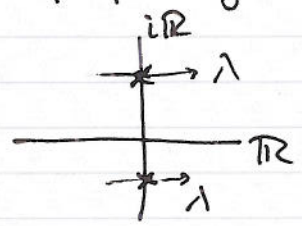
$$\Rightarrow \text{we are left to solve: } g(x, \lambda) := f_2(x, 0, \lambda) = \lambda + x^2$$

(1dim. problem).

Conclusion: typically in one-par families of ODEs locally the nr of equilibria changes as in a fold bifurcation.

(Note: in practice we cannot usually solve for $y(x, \lambda)$ explicitly, but we can find/derive a Taylor series expansion for y as fc of x and λ , and this usually suffices.)

We now consider what happens if the derivative of vector field at equilibrium point has a pair of purely imaginary eigenvalues



Example:

$$\frac{d}{dt} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} \lambda & -1 \\ 1 & \lambda \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

↑
evals $\lambda \pm i$ if $\lambda=0$ evals $\pm i$

if $\lambda < 0 \Rightarrow$ equilibrium $\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$ is asympt. stable

$\lambda > 0 \Rightarrow$ " " " " unstable

(indeed local flow changes considerably if λ passes through 0)

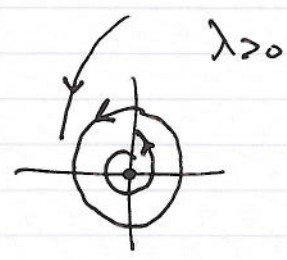
Q: What happens if the vector field is not linear?

(part of the) answer: also if the vector field is nonlinear, we see a change of stability of the equilibrium pt, if evals of Jacobian pass through $i\mathbb{R} \setminus \{0\}$

..... but is this the whole story?

Example:

$$\begin{cases} \dot{x} = \lambda x - y - x(x^2 + y^2) \\ \dot{y} = x + \lambda y - y(x^2 + y^2) \end{cases}$$



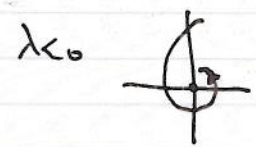
in polar coordinates
 $x = r \cos \theta$
 $y = r \sin \theta$

$$\begin{cases} \dot{r} = \lambda r - r^3 \\ \dot{\theta} = 1 \end{cases}$$

$$\dot{r} = 0 \iff r(\lambda - r^2) = 0$$

$r = 0$
 $\lambda = r^2$

"Hopf-bifurcation"



Lyapunov functions

These are useful to establish Lyapunov and/or asymptotic stability of equilibria.

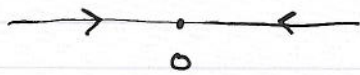
We of course already have information from derivation of vector field at eq. pt.

May be useful: (1) if eq. is not hyperbolic

(2) to obtain more info about the "region of stability" (i.e. basin of attraction in case of asymp. stable eq.)

Example

(ad 1) $\dot{x} = -x^3 = f(x) \quad x \in \mathbb{R}$



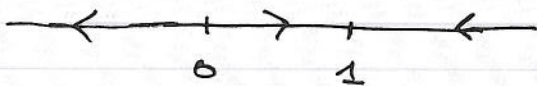
All solns $x(t)$ satisfy

$$\lim_{t \rightarrow \infty} x(t) = 0$$

$$f'(0) = 0$$

$\Rightarrow 0$ is a non-hyperbolic equilibrium pt

(ad 2) $\dot{x} = x(1-x^2) = f(x) \quad x \in \mathbb{R}$ two equilibria $x=0, x=1$



$$f'(1) = (1-2x)_{x=1} = -1 \Rightarrow x=1 \text{ is asympt stable eq. pt.}$$

from picture we see all $x(t)$ with $|x| > 0$ satisfy $\lim_{t \rightarrow \infty} x(t) = 1$

Theorem:

Let $\bar{x} \in \mathbb{R}^m$ be an equilibrium for an autonomous ODE

$$\frac{dx}{dt} = f(x) \quad x \in \mathbb{R}^m$$

and let $V: U \rightarrow \mathbb{R}$ be a differentiable function defined on some neighbourhood U of \bar{x} such that

- (1) $V(\bar{x}) = 0$ and $V(x) > 0$ if $x \neq \bar{x}$
- (2) $\frac{d}{dt} V(x(t)) \leq 0$ in U where $x(t)$ is soln of ODE

Then \bar{x} is Lyapunov stable.

Moreover, if

- (3) $\frac{d}{dt} V(x(t)) < 0$ for all $x(t) \in U \setminus \{\bar{x}\}$

then \bar{x} is ~~Lyapunov~~ asymptotically stable.

The above defined fn V is called a Lyapunov function

WARNING: there are no general methods to derive Lyapunov function for Lyapunov/asympt stable equilibria.

HOPE: there are classes of problems where there is a natural choice of Lyapunov function

(physically motivated)

e.g. mechanical systems (possibly with dissipation / friction)
have natural candidate $V = E$ (energy)

if $\frac{dE}{dt} = 0$ then local maxima and minima of energy E are Lyapunov stable equilibria.

for mechanical system with dissipation we have $\frac{dE}{dt} \leq 0$

==

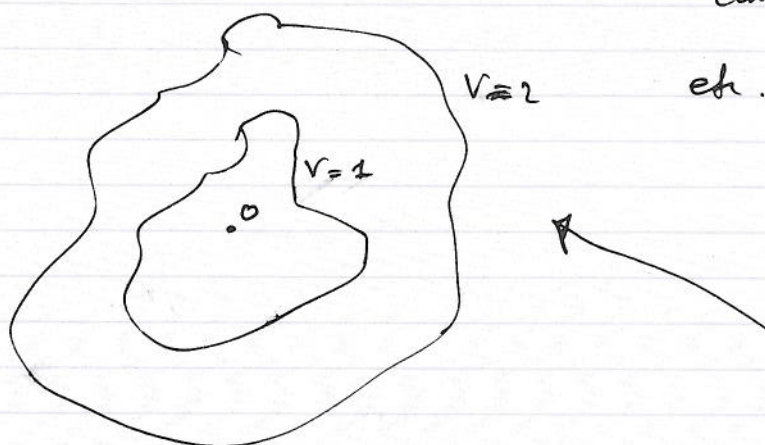
the picture (before we prove this theorem)

consider example in the plane \mathbb{R}^2 and consider $V: \mathbb{R}^2 \rightarrow \mathbb{R}$

$$V(0) = 0, \quad V(x) > 0 \quad \forall x \neq 0$$

natural to draw contour plot of V (curves at which V is const.)

$$\frac{d}{dt} V(x(t)) \leq 0$$



two scenarios: (1) along every soln V is strictly decreasing as fn of time $\Rightarrow \lim_{t \rightarrow \infty} x(t) = 0$

(2) the soln gets "stuck" on some level set of V so that $\lim_{t \rightarrow \infty} V(x(t)) = C > 0$

But in any case, ~~if~~ we are Lyapunov stable because

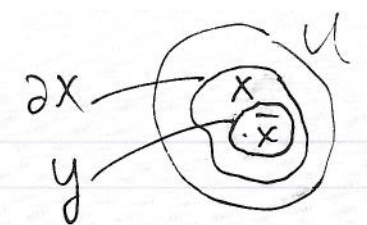
if ~~if~~ $V(x(0)) = V_0$ then we can choose a small

neighbourhood of $\bar{x} = 0$ such that in this neighbourhood

$V < V_0$ (by continuity of V) \Rightarrow any initial condition in this ball cannot intersect the level set ~~if~~ $V_0 = V(x)$

If $\frac{d}{dt} V(x(t)) < 0$ then we cannot have $\lim_{t \rightarrow \infty} V(x(t)) = C > 0$

since this would imply that $\exists y$ (accumulated by $x(t)$) such that $\frac{d}{dt} V(x(t)) = 0$ if $x(0) = y$.



Proof:

Let X be some nbh of \bar{x} , with $X \subset U$. ^{equilibrium}

Consider the function V on ∂X , the boundary of X

Since ∂X is closed and bounded and V is continuous in U

then the function V attains some minimum $V_{\min}(\partial X) = \inf_{x \in \partial X} V(x)$ somewhere on ∂X

Now consider a smaller nbh Y of \bar{x} such that

the maximum value of V on Y $V_{\max}(\partial Y) = \sup_{x \in \partial Y} V(x)$

is strictly smaller than $V_{\min}(\partial X)$. The existence ^{of Y} is guaranteed by continuity of V .

Hence, since $\frac{d}{dt} V(x(t)) \leq 0$ V cannot increase along solution curves \Rightarrow every sol. with initial condition ^{at t_0} inside Y stays inside X for all positive time.

This construction holds for all nbhs X of $\bar{x} \Rightarrow$

\bar{x} is Lyapunov stable.

Suppose, in addition, that $\frac{d}{dt} V(x(t)) < 0$. We derive a contradiction ~~to prove~~ to prove that $\lim_{t \rightarrow \infty} x(t) = \bar{x} \forall x(t_0)$ in U with $V(x(t_0)) < V_{\min}(\partial U)$.

Suppose $\lim_{t \rightarrow \infty} V(x(t)) = C > 0$ (i.e. $\lim_{t \rightarrow \infty} x(t) \neq \bar{x}$),

then there is a point y with $V(y) = C$ that is

accumulated by $x(t)$, i.e. \exists increasing sequence t_n ($n \rightarrow \infty$)

such that $\lim_{n \rightarrow \infty} x(t_n) = y \Rightarrow \frac{d}{dt} V(x(t)) = 0$ if $x(t) = y$.

exercise!