

UNIVERSITY OF LONDON  
BSc and MSci EXAMINATIONS (MATHEMATICS)  
May-June 2006

This paper is also taken for the relevant examination for the Associateship.

**M2A3**  
**One-Dimensional Fluid Mechanics**

Date: Friday, 19th May 2006      Time: 10 am – 12 noon

Credit will be given for all questions attempted but extra credit will be given for complete or nearly complete answers.

Calculators may not be used.

1. (i) Give a brief outline of the Eulerian and Lagrangian descriptions of unsteady one-dimensional fluid motion and explain how the material derivative  $D/Dt$  is defined.
- (ii) Fluid flows through a pipe with constant rectangular cross-sectional area  $A$ . The fluid velocity at a location  $x$  along the pipe is denoted by  $u(x, t)$  and the fluid density by  $\rho(x, t)$ . By concentrating on the fluid motion in a fixed region  $a \leq x \leq b$ , use the concepts of mass and momentum flux to derive the governing equations

$$\begin{aligned}\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) &= 0, \\ \frac{\partial}{\partial t}(\rho u) + \frac{\partial}{\partial x}(\rho u^2) &= -\frac{\partial p}{\partial x}.\end{aligned}$$

You may assume that the only force acting on the fluid is due to the pressure  $p(x, t)$ . Show that the second of these equations can be written in the alternative form

$$\rho \frac{Du}{Dt} = -\frac{\partial p}{\partial x}.$$

- (iii) Explain how the above equations are modified if, rather than through an enclosed pipe, the flow is along an open horizontal channel of rectangular cross-section with constant width, but variable depth  $h(x, t)$ , and the fluid is of constant density.
2. In an elastic-walled pipe of circular cross-section, the radius  $r(x, t)$  and fluid velocity  $u(x, t)$  are governed by the equations

$$\frac{\partial}{\partial t}(r^2) + \frac{\partial}{\partial x}(r^2 u) = 0, \quad \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = -\frac{2}{\alpha^2} \frac{\partial r}{\partial x}, \quad (1)$$

where  $\alpha$  is a constant.

Show that small perturbations to the rest state  $u = 0, r = r_0$  ( $r_0$  constant) propagate with speed  $c_0 = r_0^{1/2}/\alpha$ .

Now consider the steady version of (1), and show that upon integration  $r$  satisfies:

$$Q^2 = \left(\frac{2}{\alpha}\right)^2 r^4 (E - r),$$

where  $Q$  and  $E$  are constants. Interpret the constant  $Q$  physically.

Show that  $Q$  is a maximum ( $Q_{\max}$ , say) when  $r = 4E/5$ .

By sketching  $Q^2$  versus  $r$ , show graphically that, for a given value of  $Q$  ( $< Q_{\max}$ ) there are two possible values of  $r$ . What are the corresponding values of  $u$ ?

By defining an appropriate Froude number, show that one of these flows is subcritical and the other supercritical.

3. The one-dimensional flow of a river of depth  $h(x, t)$  can be modelled by the nonlinear equations

$$\frac{\partial h}{\partial t} + \frac{\partial}{\partial x}(hu) = 0,$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = -g \frac{\partial h}{\partial x} \cos \alpha + g \sin \alpha - D_0 \frac{u^2}{h}.$$

Here,  $\alpha$  is the slope of the river and the constant  $D_0$  is a drag coefficient.

- (i) Show that there is a uniform solution with  $u = u_0$ ,  $h = h_0$  and  $u_0^2 = (g \sin \alpha / D_0) h_0$ .  
(ii) Seek small perturbations to this basic state by writing

$$u = u_0 + \tilde{u}(x, t), \quad h = h_0 + \tilde{h}(x, t),$$

and show that  $\tilde{u}$  and  $\tilde{h}$  satisfy the linearised equations

$$\left( \frac{\partial}{\partial t} + u_0 \frac{\partial}{\partial x} \right) \tilde{h} = -h_0 \frac{\partial \tilde{u}}{\partial x},$$

$$\left( \frac{\partial}{\partial t} + u_0 \frac{\partial}{\partial x} \right) \tilde{u} = -g \frac{\partial \tilde{h}}{\partial x} \cos \alpha - D_0 \frac{u_0^2}{h_0} \left( \frac{2\tilde{u}}{u_0} - \frac{\tilde{h}}{h_0} \right).$$

- (iii) Show that, upon elimination of  $\tilde{u}$ , the perturbation to the depth satisfies

$$\left( \frac{\partial}{\partial t} + u_0 \frac{\partial}{\partial x} \right)^2 \tilde{h} = c_0^2 \frac{\partial^2 \tilde{h}}{\partial x^2} - g \sin \alpha \left( \frac{2}{u_0} \frac{\partial \tilde{h}}{\partial t} + 3 \frac{\partial \tilde{h}}{\partial x} \right),$$

where  $c_0^2 = gh_0 \cos \alpha$ .

- (iv) Seek wave-like solutions for  $\tilde{h}$  proportional to  $\exp[ik(x - \tilde{c}t)]$ , with  $k$  real, and derive the dispersion relation

$$\left( 3 - \frac{2\tilde{c}}{u_0} \right) g \sin \alpha + ik \left( (u_0 - \tilde{c})^2 - c_0^2 \right) = 0.$$

Consider a solution for  $\tilde{c}$ , valid for small  $k$ , of the form

$$\tilde{c} = c_1 + kc_2 + \dots$$

Show that  $c_1 = 3u_0/2$  and find  $c_2$ .

Deduce that long wavelength disturbances on the river will grow exponentially in time provided  $4D_0 < \tan \alpha$ .

4. Constant-density fluid of depth  $h$  flows along an open channel. The governing equations are

$$\frac{\partial(h^2)}{\partial t} + \frac{\partial}{\partial x}(h^2u) = 0, \quad \frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} = -g\frac{\partial h}{\partial x}.$$

- (i) Show that these equations may be written in the alternative form

$$\left(\frac{\partial}{\partial t} + (u + c)\frac{\partial}{\partial x}\right)(u + 4c) = 0, \quad \left(\frac{\partial}{\partial t} + (u - c)\frac{\partial}{\partial x}\right)(u - 4c) = 0, \quad (1)$$

where  $c^2 = \frac{1}{2}gh$ .

- (ii) Suppose that  $u = 0$  when  $h = h_0$  (constant). Deduce that equations (1) admit a solution with  $(u - 4c)$  constant everywhere, and  $u$  constant along the straight-line characteristics

$$\frac{dx}{dt} = \frac{5}{4}u + c_0, \quad c_0^2 = \frac{1}{2}gh_0.$$

The channel is bounded at  $x = 0$  by a wall, with the fluid occupying the region  $x > 0$ . Initially the fluid is at rest and has depth  $h_0$ . From time  $t = 0$  onwards the wall moves in the negative  $x$ -direction with speed  $\alpha t$  ( $\alpha$  constant).

- (iii) Assuming a solution of the form given in (ii), deduce that in the region  $x \geq c_0t$ , the flow remains in its initial state.
- (iv) By considering a characteristic originating from the wall at time  $\tau$ , show that in the region  $-\frac{1}{2}\alpha t^2 \leq x \leq c_0t$  the solution for  $u$  is

$$u(x, t) = -\alpha\tau(x, t),$$

with  $\tau$  given in terms of  $x$  and  $t$  by

$$x + \frac{1}{2}\alpha\tau^2 = \left(c_0 - \frac{5}{4}\alpha\tau\right)(t - \tau).$$

Find the corresponding solution for  $h(x, t)$ .

5. An isentropic gas evolves according to the equations

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) = 0, \quad \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = -\rho \frac{\partial \rho}{\partial x}, \quad (1)$$

in the usual notation.

(i) Show that small perturbations  $\tilde{u}(x, t), \tilde{\rho}(x, t)$  to the uniform state  $u = u_0, \rho = \rho_0$  satisfy a convected wave equation, and deduce that the wave disturbances propagate at speeds  $u_0 \pm \rho_0$ .

(ii) Show that the nonlinear equations (1) can be rewritten in the form

$$\left( \frac{\partial}{\partial t} + (u + \rho) \frac{\partial}{\partial x} \right) (u + \rho) = 0, \quad \left( \frac{\partial}{\partial t} + (u - \rho) \frac{\partial}{\partial x} \right) (u - \rho) = 0.$$

Deduce that  $(u + \rho), (u - \rho)$  are constant along straight-line characteristics, and write down the equations of these characteristics.

(iii) Hence or otherwise deduce that the general solution for  $u$  can be written as

$$u = F(x - (u + \rho)t) + G(x - (u - \rho)t),$$

where  $F, G$  are arbitrary functions. Find the corresponding solution for  $\rho$ . What is the relationship between  $u$  and  $\rho$  that ensures that waves propagate in both the positive and negative  $x$ -directions?