## PAPER 71

## STRUCTURE AND EVOLUTION OF STARS

Attempt THREE questions.
There are FOUR questions in total. The questions carry equal weight.

The symbols used in these questions have the meanings they were given in the lectures.
Candidates are reminded of the equations of stellar structure in the form:

$$
\begin{array}{cc}
\frac{d P}{d r}=-\frac{G m \rho}{r^{2}} & \frac{d m}{d r}=4 \pi r^{2} \rho \\
\frac{d T}{d r}=-\frac{3 \kappa \rho L_{r}}{16 \pi a c r^{2} T^{3}} & \frac{d L_{r}}{d r}=4 \pi r^{2} \rho \epsilon \\
P=\frac{\mathcal{R} \rho T}{\mu}+\frac{a T^{4}}{3} & \text { with }
\end{array} 1 / \mu=2 X+3 Y / 4+Z / 2
$$

STATIONERY REQUIREMENTS
Cover sheet
Treasury Tag
Script paper

SPECIAL REQUIREMENTS
None

You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.

1 A cluster of massive stars is contracting towards the main sequence with the only energy source being due to gravitational contraction. The stellar material is an ideal gas with $\gamma=5 / 3$ with radiation pressure being negligible, energy transport is by radiation and the opacity $\kappa=\kappa_{0}$, where $\kappa_{0}$ is constant.

Show that during this evolution the effective energy production rate per unit mass is given by

$$
\epsilon=-\frac{3}{2 \rho} \frac{\partial P}{\partial t}+\frac{5 P}{2 \rho^{2}} \frac{\partial \rho}{\partial t},
$$

with $t$ being the time and the derivative being taken at constant $m$.
A set of dimensionless variables are defined through $x=r / R, q=m / M, l=L_{r} / L$, $b=\left(4 \pi \rho R^{3}\right) / M$ and $p=\left(4 \pi R^{4} P\right) /\left(G M^{2}\right)$, with $q, l, b, p$ being functions only of $x$. The radius $R$ and luminosity $L$ are functions only of time.

Show that in terms of these variables, the equations of stellar structure for a contracting star take the form

$$
\begin{aligned}
\frac{d p}{d x} & =-\frac{b q}{x^{2}}, \quad \frac{d q}{d x}=x^{2} b, \\
\frac{d}{d x}\left(\frac{p}{b}\right) & =-D \frac{b^{4} l}{x^{2} p^{3}}, \quad \frac{d l}{d x}=E x^{2} p,
\end{aligned}
$$

where

$$
\begin{gathered}
D=\frac{3 \kappa_{0} \mathcal{R}^{4} L}{64 \pi^{2} a c \mu^{4} G^{4} M^{3}} \quad \text { and } \\
E=-\frac{3 G M^{2}}{2 R^{2} L}\left(\frac{d R}{d t}\right) .
\end{gathered}
$$

Hence deduce that the luminosity is $\propto M^{3}$ and constant during the evolution of a particular star. If the evolution commences at $t=0$ with very large radius, show that the radius is subsequently given by

$$
\frac{R L t}{G M^{2}}=\text { constant }
$$

The stars eventually reach the main sequence where the energy generation is by the CNO cycle with $\epsilon=\epsilon_{0} \rho T^{16}$. Show that they then obey the mass-radius relation

$$
R \propto M^{15 / 19}
$$

Show further that the mass of the stars in the cluster that are just reaching the main sequence at time $t$ satisfies a relation of the form

$$
M \propto t^{-19 / 34}
$$

2 Derive Schwarzschild's condition for stability to convection of a stellar radiative region consisting of an ideal gas with ratio of specific heats $\gamma=5 / 3$ in the form

$$
\frac{d P}{d r}>\frac{5 P}{3 \rho} \frac{d \rho}{d r}
$$

Show that this can be written alternatively as

$$
\frac{3 \kappa L_{r} P}{16 \pi a c G m T^{4}}<\frac{2}{5}
$$

The temperature in the atmosphere of a cool star is given as a function of the optical depth $\tau$ by

$$
T^{4}=T_{e}^{4}\left(\frac{1}{2}+\frac{3}{4} \tau\right)
$$

and the opacity is given by $\kappa=\kappa_{0} \rho T^{13}$, where $\kappa_{0}$ is constant.
Show that in the upper radiative layers

$$
P^{2}=\frac{4 \pi c G M a \mathcal{R}}{3 \kappa_{0} L \mu T_{e}^{8}}\left(4-T_{e}^{8} / T^{8}\right)
$$

and deduce that convetion sets in when $T=(13 / 20)^{1 / 8} T_{e}$.
In the lower convective region, the structure is polytropic with $P=K T^{5 / 2}$. Show that if the star is fully convective, there is a relation between the mass, radius and luminosity of the form

$$
L \propto M^{8 / 17} R^{38 / 17}
$$

3 (a) Show that the electron pressure in a helium gas in which the electrons are completely degenerate but nonrelativistic is given by

$$
P=K \rho^{5 / 3}
$$

where

$$
K=\left(\frac{3}{2 \pi}\right)^{2 / 3} \frac{h^{2}}{40 m_{e} m_{p}^{5 / 3}}
$$

with $h, m_{e}$, and $m_{p}$ being Planck's constant, the mass of the electron and the mass of a proton respectively.

Deduce that non relativistic helium white dwarfs obey the mass radius relation $R=A M^{-1 / 3}$, where $A$ is a constant.
(You may assume that for complete degeneracy, the number density of electrons, $n(p)$, with total momentum less than $p$, is given by $d n(p) / d p=8 \pi p^{2} / h^{3}, p<p_{0}$ and $d n(p) / d p=$ $0, p>p_{0}$, where $p_{o}$ is the Fermi momentum.)
(b) The core of a red giant has mass $M_{c}$ and is in a regime in which the radius $R_{c}$ does not vary with $M_{c}$.
Above the core is a hydrogen rich radiative envelope which is assumed to have negligible mass. The base of the envelope, at the core surface, coincides with the base of a thin hydrogen burning shell in which the luminosity $L$ is generated. The opacity is given by $\kappa=\kappa_{0} \rho / T^{7 / 2}$ and radiation pressure is neglected.

Assuming the envelope extends to small values of $P$ and $T$, show that in the regions above the shell

$$
P=C T^{17 / 4}
$$

where

$$
C=\left(\frac{64 \pi a c G M_{c} \mathcal{R}}{51 \kappa_{0} L \mu}\right)^{1 / 2}
$$

Show further that $T$ as a function of $r$ is given by

$$
T=\frac{4 \mu G M_{c}}{17 \mathcal{R} r} .
$$

The energy generation rate in the hydrogen rich layers is given by $\epsilon=\epsilon_{0} \rho T^{67 / 4}$. Confirm that the hydrogen burning shell is thin by showing that

$$
\epsilon\left(1.05 R_{c}\right) / \epsilon\left(R_{c}\right) \sim 1 / e
$$

and deduce that the luminosity-core mass relation is

$$
L \propto M_{c}^{97 / 8} .
$$

State very briefly what conditions $M_{c}$ and $R_{c}$ should satisfy so that this model is consistent.

4 A binary system with components of mass $M_{1}$ and $M_{2}$ is in circular orbit about the centre of mass with period $P_{\text {orb }}=2 \pi / \Omega$. Their distances from the centre of mass are $a_{1}$ and $a_{2}$ and their angular momenta about the centre of mass are $J_{1}$ and $J_{2}$ respectively. The separation $a=a_{1}+a_{2}$, and the total orbital angular momentum $J=J_{1}+J_{2}$.

Show that

$$
J=\left(\frac{M_{1} M_{2}}{M_{1}+M_{2}}\right) a^{2} \Omega=\frac{G^{2 / 3} P_{o r b}^{1 / 3} M_{1} M_{2}}{(2 \pi)^{1 / 3}\left(M_{1}+M_{2}\right)^{1 / 3}}
$$

The star of mass $M_{1}$ is transferring mass to $M_{2}$ and simultaneously losing mass to infinity through a stellar wind. The mass transfer rate is $\dot{M}_{2}=-f \dot{M}_{1}$ and the mass loss rate due to the wind is $(1-f) \dot{M}_{1}$. The wind carries away a specific angular momentum $J_{1} / M_{1}$.
By considering the conservation of angular momentum or otherwise, deduce that

$$
P_{o r b} \propto M_{1}^{-3 f} M_{2}^{-3}\left(M_{1}+M_{2}\right)^{-2} .
$$

The Roche lobe of $M_{1}$ is given by the relation

$$
R_{L}=0.46 a\left(\frac{M_{1}}{\left(M_{1}+M_{2}\right)}\right)^{1 / 3}
$$

Show that

$$
\frac{1}{R_{L}} \frac{d R_{L}}{d t}=\frac{\dot{M}_{1}}{M_{1}}\left(f\left(2 q+\frac{4 q}{3(1+q)}-2\right)+\frac{1}{3}-\frac{4 q}{3(1+q)}\right)
$$

where $q=M_{1} / M_{2}$.
The radius $R_{1}$ of $M_{1}$ is such that $R_{1} \propto M_{1}^{-n}$ and it is assumed to remain in contact with the Roche lobe. Deduce that in that case $f$ must be such that

$$
f\left(2 q+\frac{4 q}{3(1+q)}-2\right)=\frac{4 q}{3(1+q)}-n-\frac{1}{3}
$$

Comment on what happens when $q$ and $n$ are such that the above expression returns a result for $f$ outside the interval $(0,1)$.

## END OF PAPER

