

Tuesday 29 May 2007 9 to 12

MATHEMATICS (1)

Before you begin read these instructions carefully:

*You may submit answers to no more than **six** questions. All questions carry the same number of marks.*

The approximate number of marks allocated to a part of a question is indicated in the right hand margin.

*Write on **one** side of the paper only and begin each answer on a separate sheet.*

At the end of the examination:

*Each question has a number and a letter (for example, **6A**).*

*Answers must be tied up in **separate** bundles, marked **A, B or C** according to the letter affixed to each question.*

Do not join the bundles together.

For each bundle, a blue cover sheet must be completed and attached to the bundle.

*A **separate** yellow master cover sheet listing all the questions attempted **must** also be completed.*

Every cover sheet must bear your examination number and desk number.

STATIONERY REQUIREMENTS

6 blue cover sheets and treasury tags

Yellow master cover sheet

Script paper

SPECIAL REQUIREMENTS

None

<p>You may not start to read the questions printed on the subsequent pages until instructed to do so by the Invigilator.</p>

1A

For the spherical polar co-ordinates (r, θ, ϕ) defined by

$$x_1 = r \sin \theta \cos \phi, \quad x_2 = r \sin \theta \sin \phi, \quad x_3 = r \cos \theta, \quad r > 0, \quad 0 \leq \theta \leq \pi, \quad 0 \leq \phi < 2\pi$$

the gradient ∇ and the Laplacian ∇^2 are given by the following:

$$\nabla = \mathbf{e}_r \frac{\partial}{\partial r} + \frac{\mathbf{e}_\theta}{r} \frac{\partial}{\partial \theta} + \frac{\mathbf{e}_\phi}{r \sin \theta} \frac{\partial}{\partial \phi},$$

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right].$$

Let

$$\hat{\mathbf{r}} = \frac{\mathbf{r}}{r}.$$

(i) Show that

$$\nabla \times \hat{\mathbf{r}} = 0, \quad \nabla \cdot \hat{\mathbf{r}} = \frac{2}{r}.$$

[5]

(ii) Let

$$\mathbf{f}(\mathbf{r}) = R\hat{\mathbf{r}} + \hat{\mathbf{r}} \times \nabla F + (\hat{\mathbf{r}} \times \nabla G) \times \hat{\mathbf{r}}$$

where R, F, G are smooth functions of r, θ, ϕ . Show that

(a)

$$\mathbf{f} = R\mathbf{e}_r + \frac{1}{r} \left(\frac{\partial G}{\partial \theta} - \frac{1}{\sin \theta} \frac{\partial F}{\partial \phi} \right) \mathbf{e}_\theta + \frac{1}{r} \left(\frac{1}{\sin \theta} \frac{\partial G}{\partial \phi} + \frac{\partial F}{\partial \theta} \right) \mathbf{e}_\phi.$$

[5]

(b) $(\nabla \times \mathbf{f}) \cdot \mathbf{r}$ is independent of R and G .

[5]

(c)

$$(\nabla \times \mathbf{f}) \cdot \mathbf{r} = \frac{1}{r} \Delta_{\theta, \phi} F$$

for some differential operator $\Delta_{\theta, \phi}$ which should be determined.

[5]

[You may assume the identity $\nabla \times (\mathbf{f}_1 \times \mathbf{f}_2) = \mathbf{f}_1 \nabla \cdot \mathbf{f}_2 - \mathbf{f}_2 \nabla \cdot \mathbf{f}_1 + (\mathbf{f}_2 \cdot \nabla) \mathbf{f}_1 - (\mathbf{f}_1 \cdot \nabla) \mathbf{f}_2$.]

2A The heat flow along a thin circular wire of length $2L$ can be approximated by the heat equation

$$\frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2}, \quad \kappa > 0, \quad -L < x < L$$

with the following boundary conditions

$$u(-L, t) = u(L, t), \quad \frac{\partial u}{\partial x}(-L, t) = \frac{\partial u}{\partial x}(L, t).$$

(i) Use separation of variables to express the solution u in terms of an infinite Fourier series involving appropriate integral transforms of the initial condition $u(x, 0) = u_0(x)$.

[10]

(ii) Compute the integral transforms appearing in (i) in the particular case that

$$u_0(x) = \sin x - \left(\frac{\sin L}{L}\right)x.$$

[10]

3A

(i) The cosine transform of a function $f(x)$, which has sufficient smoothness and decay as $x \rightarrow \infty$, is given by

$$\hat{f}_c(k) = 2 \int_0^{\infty} f(x) \cos kx \, dx, \quad k > 0 .$$

By applying the Fourier transform to an even function, show that the inverse cosine transform is given by

$$f(x) = \frac{1}{\pi} \int_0^{\infty} \hat{f}_c(k) \cos kx \, dk, \quad x > 0 .$$

[7]

(ii) Let $f(x, t)$ satisfy the partial differential equation

$$i \frac{\partial f}{\partial t} + \frac{\partial^2 f}{\partial x^2} = 0, \quad -\infty < x < \infty, \quad t > 0 ,$$

and the initial condition

$$f(x, 0) = f_0(x), \quad -\infty < x < \infty ,$$

where the function $f_0(x)$ has sufficient smoothness and decay as $|x| \rightarrow \infty$. Use the Fourier transform to show that $f(x, t)$ can be written as

$$f(x, t) = \frac{c}{2\pi\sqrt{t}} \int_{-\infty}^{\infty} f_0(\xi) e^{\frac{i(x-\xi)^2}{4t}} \, d\xi ,$$

where

$$c = \int_{-\infty}^{\infty} e^{-it^2} \, dt .$$

[13]

4C

Define the trace of a square matrix and show that $\text{Tr}(AB) = \text{Tr}(BA)$. Deduce that there are no $n \times n$ matrices A, B such that

$$AB - BA = I,$$

where I is the identity matrix. [6]

Let A, B be real $n \times n$ matrices such that the complex matrix $C = A + iB$ is invertible. By considering $\det(A + \lambda B)$ as a function of λ , show that the matrix $A + \lambda B$ is invertible for some real number λ . [6]

Deduce that if two real matrices P, Q are related by a complex similarity transformation $P = RQR^{-1}$, where R is a complex matrix, then they are also related by a real similarity transformation. [8]

[Hint: for the last part rearrange the similarity relation and consider its real and imaginary parts.]

5C

Define the standard inner product in a complex vector space \mathbb{C}^n , and prove the Cauchy–Schwarz inequality. [8]

Show that if U is a unitary matrix, then $|U\mathbf{a}| = |\mathbf{a}|$ for all vectors \mathbf{a} . Hence, find a constraint for the eigenvalues of U . [4]

Given a one parameter family of unitary matrices

$$U(t) = I + tA + O(t^2),$$

where t is real, show that the eigenvalues of the matrix A are purely imaginary. [8]

6A

(i) Let $y(z)$ satisfy

$$2z(1-z)y'' + (1+z)y' - y = 0 .$$

Find the indicial equation associated with the singular point $z = 1$.

[4]

(ii) Let $\psi(x)$ satisfy the Schrödinger equation

$$\psi'' + (2\lambda + 1 - x^2)\psi = 0 .$$

Use the transformation $\psi(x) = e^{-\frac{x^2}{2}}y(x)$ to obtain the equation satisfied by $y(x)$, which is called the Hermite equation.

[8]

Construct two linearly independent series solutions of the Hermite equation in the neighbourhood of $x = 0$. Give the radius of convergence of the series obtained and construct the first three terms of these series. Find the particular values of λ for which there exist polynomial solutions. Find such solutions up to terms including x^3 .

[8]

7A

(i) Let $y_n(x)$ satisfy

$$\frac{d}{dx}\left(g(x)\frac{dy_n(x)}{dx}\right) + h(x)y_n(x) + \lambda_n w(x)y_n(x) = 0, \quad a < x < b,$$

$$\alpha_1 \frac{dy_n}{dx}(a) + \alpha_2 y_n(a) = 0, \quad \beta_1 \frac{dy_n}{dx}(b) + \beta_2 y_n(b) = 0,$$

where α_1, α_2 are not both zero and β_1, β_2 are not both zero. Show that if $\lambda_m \neq \lambda_n$ then y_n and y_m satisfy the following orthogonality relation

$$\int_a^b w(x)y_n(x)y_m(x) dx = 0,$$

for $m \neq n$.

[10]

(ii) Let

$$\frac{d}{dx}\left(x\frac{dy_n(x)}{dx}\right) + \frac{\lambda_n}{x}y_n(x) = 0, \quad 1 < x < b,$$

$$y_n(1) = y_n(b) = 0,$$

where $\lambda_n \neq 0$. Use the change of variables $\xi = \lambda_n \ln x$ to compute λ_n and $y_n(x)$.

[10]

8A Use an appropriate Green's function to solve the following differential equation:

$$\frac{d^2 u}{dx^2} - u = e^x, \quad 0 < x < 1,$$

$$u(0) = u(1) = 0.$$

[20]

[You may use the identity $\sinh a \cosh b - \cosh a \sinh b = \sinh(a - b)$.]

9B

Show that functions $y(x)$ which make stationary the functional

$$F[y] = \int_b^a f(x, y, y') dx ,$$

where $y' = \frac{dy}{dx}$, and $y(a), y(b)$ are fixed, satisfy Euler's equation

$$\frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = \frac{\partial f}{\partial y} .$$

[10]

The motion of a particle in a plane is constrained by the Lagrangian

$$L = \frac{\dot{r}^2}{\left(1 - \frac{2m}{r}\right)} + r^2 \dot{\theta}^2 + \frac{\mu}{\left(1 - \frac{2m}{r}\right)} ,$$

where m, μ are positive constants and $r(t), \theta(t)$ are generalized co-ordinates with $r > 2m$, $0 \leq \theta < 2\pi$. By setting

$$J = r^2 \dot{\theta} ,$$

show that the Euler-Lagrange equations imply that J is constant.

By computing the Euler-Lagrange equation for r , show that solutions with $r = R$ for constant $R > 2m$ are only possible when

$$J^2 = \frac{m\mu R}{\left(1 - \frac{2m}{R}\right)^2} .$$

[10]

10B

The Sturm-Liouville equation is given by

$$-\frac{d}{dx}(p(x)y') + q(x)y = \lambda w(x)y, \quad a < x < b, \quad (\star)$$

where λ is constant and $p(x) > 0$, $q(x) > 0$ and $w(x) > 0$. Show that solutions of this equation which satisfy the boundary condition

$$p(b)y(b)y'(b) - p(a)y(a)y'(a) = 0$$

correspond to functions $y(x)$ for which the quotient

$$\Lambda[y] = \frac{F[y]}{G[y]}$$

is stationary, where

$$F[y] = \int_a^b (py'^2 + qy^2)dx, \quad G[y] = \int_a^b \omega y^2 dx.$$

Show furthermore that the eigenvalues λ of this Sturm-Liouville problem are given by the values of $\Lambda[y]$, where y satisfies (\star) .

[10]

Suppose $y(x)$ satisfies the second order differential equation

$$-y'' - \frac{(n-1)}{x}y' + x^2y = \lambda y, \quad x > 0,$$

where n is a fixed positive integer and λ is constant. By writing this equation in the form of a Sturm-Liouville equation with $p(x) = w(x) = x^{n-1}$ and $q(x) = x^{n+1}$, use the Rayleigh-Ritz method with a trial function of the form e^{ax^2} for $a < 0$ to find an estimate for the lowest eigenvalue of the Sturm-Liouville problem.

[10]

[For this question, you may assume the Euler equation without proof.]

END OF PAPER